SN 2008in – BRIDGING THE GAP BETWEEN NORMAL AND FAINT SUPERNOVAE OF TYPE IIP

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ABSTRACT

We present optical photometric and low-resolution spectroscopic observations of the Type II plateau supernova (SN) 2008in, which occurred in the outskirts of the nearly face-on spiral galaxy M 61. Photometric data in the X-rays, ultraviolet and near-infrared bands have been used to characterize this event. The SN field was imaged with the ROTSE-IIIb optical telescope about seven days before the explosion. This allowed us to constrain the epoch of the shock breakout to JD = 2454825.6. The duration of the plateau phase, as derived from the photometric monitoring, was ~ 98 days. The spectra of SN 2008in show a striking resemblance to those of the archetypal low-luminosity IIP SNe 1997D and 1999br. A comparison of ejecta kinematics of SN 2008in with the hydrodynamical simulations of Type IIP SNe by Dessart et al. (2010) indicates that it is a less energetic event ($\sim 5 \times 10^{50}$ erg). However, the light curve indicates that the production of radioactive ⁵⁶Ni is significantly higher than that in the low-luminosity SNe. Adopting an interstellar absorption along the SN direction of $A_V \sim 0.3$ mag and a distance of 13.2 Mpc, we estimated a synthesized ⁵⁶Ni mass of $\sim 0.015 M_{\odot}$. Employing semi-analytical formulae (Litvinova & Nadezhin 1985), we derived a pre-SN radius of ~ $126R_{\odot}$, an explosion energy of $\sim 5.4 \times 10^{50}$ erg and a total ejected mass of $\sim 16.7 M_{\odot}$. The latter indicates that the zero age main-sequence mass of the progenitor did not exceed $20M_{\odot}$. Considering the above properties of SN 2008in and its occurrence in a region of sub-solar metallicity ($[O/H] \sim 8.44 \text{ dex}$), it is unlikely that fall-back of the ejecta onto a newly formed black hole occurred in SN 2008in. We therefore favor a low-energy explosion scenario of a relatively compact, moderate-mass progenitor star that generates a neutron star.

Subject headings: supernovae: general – supernovae: individual (2008in)

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1. INTRODUCTION

Core-collapse Type II supernovae (SNe) mark the violent death of stars with main-sequence masses greater than 8 M_{\odot} and as is indicated by the presence of hydrogen lines in their optical spectra, they originate from a progenitor star with a significant amount of hydrogen still intact (Eldridge & Tout 2004). Of special interest are the Type II plateau SNe which are characterized by a 'plateau' in their optical light curve and are more common, constituting about 75% of all Type II SNe (Smith et al. 2011). The IIP SNe show a wide range of plateau luminosities, plateau durations, expansion velocities and nickel masses (Hamuy 2003) and these observational properties are connected with the explosion mechanism as well as the physical properties of the progenitor star such as ejected mass, explosion energy and pre-SN radius (Nadyozhin 2003; Smartt et al. 2009). The IIP SNe are thought to result from progenitor masses in the range $8-25 M_{\odot}$ (Heger et al. 2003) with an extended hydrogen envelope necessary to maintain the plateau phase. A detailed study of optical light curves and spectra of only a few nearby IIP SNe has been done so far and there exists a discrepancy in estimating the mass of their progenitors, e.g. for the three well studied events (namely 1999em, 2005cs and 2004et), the determination of progenitor mass from the hydrodynamical modeling of their light curve is found to be higher than that estimated from pre-SN imaging (Utrobin et al. 2010; Bersten et al. 2011).

Recently, a number of 'low luminosity' Type IIP events have been discovered viz. SNe 1999br, 1999eu, 1994N, 2001dc, 2005cs (Pastorello et al. 2004, 2009), 2008bk (Van Dyk et al. 2010) and 2009md (Fraser et al. 2010). These events have explosion energy ($\sim 10^{50}$ erg) and ejected ⁵⁶Ni-mass $(2-7 \times 10^{-3} M_{\odot})$, both lower by a factor of 10 than normal, and low expansion velocity ~ 1000 ${\rm km\,s^{-1}}({\rm Pastorello}~{\rm et}~{\rm al.}~2009$ and references therein). The low-luminosity IIP SNe are debated because of the unknown nature of their progenitors. The first reported faint SN was SN 1997D (Turatto et al. 1998; Benetti et al. 2001), and the observed properties of its light curves and spectra were explained in terms of significant fallback of ejected material on a newly formed black hole (BH), created through the core collapse of a massive progenitor $(M \gtrsim 20 M_{\odot}, \text{Zampieri et al. 1998, 2003})$. Alternatively, SN 1997D was interpreted as the explosion of a less massive progenitor $(8 - 12M_{\odot})$, Chugai & Utrobin 2000), close in mass to the lower limit for stars that can undergo core-collapse. Heger et al. (2003) suggested that low-luminosity Type IIP events are electron capture SNe produced by low-mass progenitors giving rise to ONeMg cores. This is further supported through investigations of pre-explosion images (Maund & Smartt 2005; Maund et al. 2005), though Eldridge et al. (2007)have ruled out the possibility of such a mechanism for the low-luminosity SN 2005cs. According to the formalism of Heger et al. (2003) and Eldridge & Tout (2004), no star having initial mass less than $22M_{\odot}$ can form a BH, which can quench the ejected material and produce a low luminosity SN. Stars with masses above $25M_{\odot}$, formed in metal-poor or slightly sub-solar metallicity regions can produce low-luminosity, BH-forming Type IIP SNe. Type IIL/b events can be produced through this process from stars having masses $> 25 M_{\odot}$ and generated in regions with solar (or super-solar) metallicity. So, the metallicity information at the SN location and the estimation of the initial mass are essential to constrain the triggering mechanisms of these explosions.

SN 2008in was discovered in the nearby galaxy M 61 (NGC 4303). The first unfiltered CCD images of SN 2008in were taken by Koichi Itagaki on 2008 December 26.79 (all times in UT hereafter) and 27.69 at a magnitude of 14.9. Independent observations of this event by K. Kadota showed the transient at an unfiltered mag of 15.1, In addition W. Wells recorded the SN on 2008 December 28.46 at V and R band magnitudes of 14.3 and 13.2, respectively (Nakano et al. 2008). Low and midresolution spectroscopic observations indicate an early discovery for SN 2008in (within 1-2 weeks after corecollapse). The spectra showed highly blueshifted H_{α} and H_{β} absorptions (by ~ 9000 km s⁻¹) with weaker emission components (Chakraborti et al. 2008; Foley 2008; Stritzinger 2008). The presence of prominent P-Cygni profiles of Balmer lines leads to its classification as a Type II SN.

The broadband light curve and the initial spectral evolution of SN 2008in were similar to those of normal Type IIP SNe. However, from mid-plateau, the SN started to show a few spectral features (like H α) which are similar to under-luminous events. SN 2008in was also observed in the radio with the Very Large Array (VLA) on 2008 December 31.40 UT in two frequency bands at 8.4601 and 22.4601 GHz (Stockdale et al. 2008). Observation for the second epoch was further reported by Stockdale et al. (2009) on 2009 January 27. Interacting Type II SNe, like Type IIn events and a few Type IIP events (e.g. SN 2004et) are supposed to be strong sources of radio emission (for review see Weiler et al. 2002). However, both VLA observations produced null results for this proximate event.

In this paper we present optical and near-infrared photometric and optical spectroscopic observations of SN 2008in. The photometric data cover a time span of about 410 days since the discovery. The Swift/XRT (Burrows et al. 2005) and Swift/UVOT (Roming et al. 2005) data covering a time span of 60 days are also presented. In Sections 3 and 4 we study the the photometric and spectroscopic evolutions respectively. In Sections 5 and 6 we cover the estimates of distance, reddening, intrinsic color and bolometric fluxes. The main physical parameters of the explosion and the mass of the progenitor are derived in Section 7 while a comparison of its properties with other SNe is given in Section 8. The conclusion of the paper is given in the last section. The epoch of explosion JD = 2454825.6 (§3) is considered throughout the paper and the times of pre-/postexplosion are referred with -/+ signs respectively.

2. OBSERVATION AND DATA REDUCTION

2.1. Photometric Observation

The prompt follow-up of the event was carried out by the ground-based ROTSE-IIIb telescope¹⁸ having sensitivity in the wavelength region from 0.35 to 1.0 μ m with

¹⁸ The Robotic Optical Transient Search Experiment (ROTSE-III) is a set of four 45 cm fully robotic optical telescopes, installed at Siding Spring Observatory, Australia (ROTSE-IIIa), Mc-

a peak around 0.6 μ m (Quimby et al. 2007). The SN was first detected in the ROTSE-IIIb images on 2008 December 24.45 and it was monitored at 58 phases until +115d. The initial detections of the SN evaded the automated pipeline identification due to poor image quality and low signal-to-noise ratio (S/N). The data reduction was performed afresh on all the available ROTSE data. In order to remove the contamination of the true SN-flux from the underlying galaxy, a pre-SN galaxy template was constructed from images taken in early 2008 and each SN frame was reduced using the galaxy-template subtraction scheme developed by Alard (2000). The point spread function (PSF) photometry was performed at the SN location in the galaxy-template subtracted images. The unfiltered instrumental magnitudes were calibrated using the USNO B1.0 *R*-band magnitudes of about 15 isolated stars. The light curve thus produced was found to be 0.15 mag off from the Cousins R band light curve produced by multi-band observation carried out at ARIES (described below) and the ROTSE magnitudes were scaled accordingly.

The SN 2008in was also monitored with the Ultraviolet Optical Telescope (UVOT) on board Swift from +5dto +60d. The UVOT filters uvw2, uvm2, uvw1, u, b, and v have their effective wavelength at 2030, 2231, 2634, 3501, 4329, and 5402 Å, respectively (Poole et al. 2008). The UVOT data reduction was performed following the prescriptions of Brown (2009). A 5" aperture is used to measure the counts for the coincidence loss correction whereas a 3'' aperture was used for the photometry. For the filters uvw2, uvm2, and uvw1, the last epoch (obtained about 60 days after explosion when the UV flux was very weak) data were used to subtract the galaxy light, while in the optical the SN remains much brighter than the underlying light so contamination was negligible. An aperture correction (based on an average PSF in *Swift* CALDB) as well as zeropoints from Poole et al. (2008) was applied to put the magnitudes on the UVOT photometric system. The UVOT magnitudes are listed in Table 1.

From +5d to +416d, the follow-up of SN 2008in in optical broadband Johnson UBV and Cousins RI was performed with the 1-m Sampurnanand Telescope (ST) at ARIES, Nainital¹⁹. An identification chart showing the field of the galaxy M61 along with the locations of SN 2008in as well as the local standards is presented in Figure 1. The photometry is performed using standard tasks of IRAF ²⁰ and DAOPHOT ²¹ as described in the paper

Donald Observatory, Texas (ROTSE-IIIb), H.E.S.S. site, Namibia (ROTSE-IIIc) and TUBITAK National Observatory, Turkey (ROTSE-IIId; Akerlof et al. 2003; Yuan 2010).

 $^{19}\,$ A 2048 \times 2048 CCD camera mounted at the f/13 Cassegrain focus of the telescope, has a square pixel of 24 micron on a side, and with a plate scale of 0.38 arcsec per pixel, it covers a square area of 13' on a side in the sky. The gain and readout noise of the CCD camera are 10 electrons per analog-to-digital unit and 5.3 electrons, respectively. To improve the S/N ratio and optimize the sampling, the observations were carried out in a binned mode of 2×2 pixel.

 20 IRAF stands for Image Reduction and Analysis Facility distributed by the National Optical Astronomy Observatory which is operated by the Association of Universities for research in Astronomy, Inc. under co-operative agreement with the National Science Foundation

 $^{21}\;\; DAOPHOT$ stands for Dominion Astrophysical Observatory Photometry (Stetson 1987)

by Roy et al. (2011). Bias subtraction and flat fielding were performed on the raw frames and the cosmic removal was done using LACOSMIC routine (van Dokkum 2001). As the SN lies in the outskirts of the galaxy on a relatively faint and smooth background, the photometry at the initial phases (mostly during the plateau phase when the SN is bright) is estimated using the profile fitting method. During nebular phases, when the SN becomes faint, the true SN flux is estimated using the galaxy template subtraction method following the procedures of Roy et al. (2011). As a galaxy template, we used post-explosion (+600d) images observed on 2011 January 04 under good seeing conditions. Figure 2 shows an example of this procedure applied to a late-time V-band image of SN 2008in. In the template image of 2011 January 04, we can see a clear flux enhancement located approximately 5" away (a linear distance of \sim 343pc) from the SN position. This knot is also present in a deep image from the Sloan Digital Sky Survey (SDSS) and it is identified as a HII region in the NED catalog. The progenitor of SN 2008in may have a possible association with this star-forming region. The field of SN 2008in was calibrated using Landolt (2009) standard stars of the field SA98 observed on the same night as the SN. A sample of 10 bright and isolated non-variable stars in the field of SN 2008in was used as local standards to derive the zero points for the SN at each epoch. The location and magnitudes of these local standards are listed in Table 2. The entire time span of photometric observation is about 410 days and in Table 3, we report the UBVRIphotometry of the event.

From +4d to +116d, the SN 2008in was also observed in JHK near-infrared (NIR) bands with the 0.6-m REM/REMIR at La Silla (Zerbi et al. 2004). The object was clearly visible in the J and H passbands, while it remained undetected in K. The instrumental magnitudes were standardized using the 2MASS standards available in the field of SN 2008in. The calibrated JH magnitudes are listed in Table 4.

2.2. X-ray Observations

The Swift XRT observations were obtained simultaneously with the UVOT observations. To search for Xray emission from SN 2008in, we extracted X-ray counts from a circular region with a 10 pixel radius (23".7, corresponding to the XRT on-axis 90% encircled energy radius) centered on the optical position of the SN. The background was extracted locally from a source-free region of 40" radius to account for the detector, sky background, and the diffuse emission from the host galaxy.

No X-ray source is detected in the merged 27.1 ks XRT data obtained in photon-counting mode. The 3σ upper limit to the XRT net count rate is 7.2×10^{-4} cts s⁻¹, corresponding to an unabsorbed (0.2–10 keV band) X-ray flux of $f_{0.2-10} < 3.4 \times 10^{-14}$ erg cm⁻² s⁻¹ and a luminosity of $L_{0.2-10} < 7.0 \times 10^{38}$ erg s⁻¹ for an adopted thermal plasma spectrum with a temperature of kT = 10 keV (see Fransson et al. 1996 and references therein), a Galactic foreground column density of $N_{\rm H} = 1.67 \times 10^{20}$ cm⁻² (Dickey & Lockman 1990) and a distance of 13.19 Mpc (§5).

2.3. Optical Spectroscopic Observation

		TABLI	Ξ 1			
Гне	$\mathit{swift}/\mathrm{UVOT}$	PHOTOMETRIC	OBSERVATIONS	OF	SN	2008in.

UT Date (yy/mm/dd)	JD 2454000+	Phase ^a (day)	uvw2 (mag)	uvm2 (mag)	uvw1 (mag)	$u \pmod{(\mathrm{mag})}$	$b \pmod{(\mathrm{mag})}$	$v \pmod{(\mathrm{mag})}$
2008/12/30.08 2008/12/30.60 2008/12/31.96 2009/01/01.83 2009/01/03.54 2009/01/07.79 2009/01/11.15 2009/01/19.79	$\begin{array}{c} 830.58\\ 831.11\\ 832.46\\ 833.30\\ 835.05\\ 839.29\\ 842.64\\ 851.30\end{array}$	5.00 5.53 6.88 7.72 9.47 13.71 17.06 25.72	$\begin{array}{c} 14.97{\pm}0.05\\ 15.13{\pm}0.05\\ 15.53{\pm}0.05\\ 15.72{\pm}0.05\\ 16.23{\pm}0.05\\ 17.87{\pm}0.08\\ 18.81{\pm}0.11\\ \end{array}$	$\begin{array}{c} 14.65{\pm}0.04\\ 14.85{\pm}0.04\\ 15.28{\pm}0.09\\ 15.47{\pm}0.05\\ 16.07{\pm}0.05\\ 17.88{\pm}0.11\\ 18.66{\pm}0.14\\ \end{array}$	$\begin{array}{c} 14.45{\pm}0.05\\ 14.54{\pm}0.04\\ 14.84{\pm}0.05\\ 15.01{\pm}0.05\\ 15.51{\pm}0.05\\ 16.96{\pm}0.07\\ 17.80{\pm}0.09\\ 18.67{\pm}0.16\end{array}$	$\begin{array}{c} 14.12{\pm}0.04\\ 14.16{\pm}0.04\\ 14.32{\pm}0.04\\ 14.32{\pm}0.04\\ 14.52{\pm}0.04\\ 15.15{\pm}0.04\\ 15.74{\pm}0.04\\ 16.67{\pm}0.05 \end{array}$	$\begin{array}{c} 15.24{\pm}0.04\\ 15.20{\pm}0.04\\ 15.25{\pm}0.04\\ 15.30{\pm}0.04\\ 15.35{\pm}0.04\\ 15.49{\pm}0.04\\ 15.62{\pm}0.04\\ 16.00{\pm}0.04\\ \end{array}$	$\begin{array}{c} 15.20{\pm}0.04\\ 15.13{\pm}0.05\\ 15.15{\pm}0.06\\ 15.17{\pm}0.04\\ 15.25{\pm}0.04\\ 15.25{\pm}0.05\\ 15.27{\pm}0.04\\ 15.36{\pm}0.05\\ \end{array}$
2009/01/22.00 2009/01/22.73	884.53 885.23	$58.95 \\ 59.65$	_	_	_	17.92 ± 0.07 18.02 ± 0.09	16.77 ± 0.05 16.75 ± 0.05	15.69 ± 0.04 15.68 ± 0.05

^a With reference to the explosion epoch JD 2454825.6



FIG. 1.— SN 2008in in M 61. A 600 s V-band image taken at phase +32d from the 1-m ST ARIES, India and covering an area of about $10' \times 10'$ is shown. The location of SN 2008in as well as the local standard stars are marked with circles. North is up and east is to the left.

TABLE 2 The photometric magnitudes of secondary standard stars in the field of SN 2008in.

Star ID	$\left(\begin{array}{c} \alpha_{ m J2000} \ ({ m h~m~s}) \end{array} ight)$	$^{\delta_{ m J2000}}_{(\circ\prime\prime\prime)}$	$U \ (mag)$	$B \pmod{(\mathrm{mag})}$	$V \pmod{(\mathrm{mag})}$	$R \pmod{(\mathrm{mag})}$	$I \pmod{(\mathrm{mag})}$
1	$12 \ 21 \ 38$	$04 \ 30 \ 26$	$16.99 {\pm} 0.03$	$17.69 {\pm} 0.03$	$17.33 {\pm} 0.05$	$16.96 {\pm} 0.06$	$16.59 {\pm} 0.05$
2	$12 \ 21 \ 40$	$04 \ 28 \ 18$	_	$19.49 {\pm} 0.05$	$18.38 {\pm} 0.04$	$17.85 {\pm} 0.01$	$17.40 {\pm} 0.02$
3	$12 \ 21 \ 44$	$04 \ 25 \ 24$	$17.91 {\pm} 0.04$	$18.09 {\pm} 0.02$	$17.65 {\pm} 0.01$	$17.30 {\pm} 0.02$	$16.94 {\pm} 0.01$
4	$12 \ 21 \ 48$	$04 \ 32 \ 27$	$18.38 {\pm} 0.11$	$17.61 {\pm} 0.01$	$16.02 {\pm} 0.01$	$15.10 {\pm} 0.01$	$14.34 {\pm} 0.01$
5	$12 \ 21 \ 50$	$04 \ 34 \ 20$	$19.10 {\pm} 0.21$	18.45 ± 0.04	17.25 ± 0.01	$16.55 {\pm} 0.01$	$15.99 {\pm} 0.01$
6	$12 \ 21 \ 57$	$04 \ 34 \ 38$	$16.06 {\pm} 0.01$	15.72 ± 0.00	$14.88 {\pm} 0.00$	$14.36 {\pm} 0.01$	$13.94 {\pm} 0.00$
7	$12 \ 22 \ 04$	$04 \ 32 \ 30$	$16.66 {\pm} 0.02$	$16.30 {\pm} 0.00$	$15.48 {\pm} 0.01$	15.00 ± 0.01	$14.61 {\pm} 0.01$
8	$12 \ 22 \ 05$	$04 \ 30 \ 51$	$18.05 {\pm} 0.06$	$18.17 {\pm} 0.01$	$17.60 {\pm} 0.04$	17.25 ± 0.03	$16.93 {\pm} 0.01$
9	$12 \ 21 \ 43$	$04 \ 31 \ 54$	$19.56 {\pm} 0.41$	$19.20 {\pm} 0.05$	17.55 ± 0.01	$16.53 {\pm} 0.01$	$15.51 {\pm} 0.01$
10	$12 \ 21 \ 33$	$04 \ 29 \ 29$	$19.11 {\pm} 0.43$	$18.96 {\pm} 0.02$	$18.55 {\pm} 0.02$	$18.21 {\pm} 0.05$	$17.89 {\pm} 0.02$

Note. — Errors in magnitude represent RMS scatter of the night-to-night repeatability over entire period of SN monitoring.



FIG. 2.— Measurement procedure of the true SN flux in the nebular phase. Each image shows $76'' \times 76''$ around SN location. North is up and east is to the left. The leftmost panel shows V-band image observed on 2010 January 20 with a SN in it. The middle panel shows the template image taken on 2011 January 04, with no SN in it. A small flux enhancement marked with a cross (×) symbol is probably due to a star forming region. The rightmost panel shows the subtracted image, where only SN is present.

TABLE 3 The UBVRI photometry of SN 2008in.

UT Date	JD	Phase ^a	U	В	V	R	Ι	Seeing ^c
(yy/mm/dd)	2454000 +	(day)	(mag)	(mag)	(mag)	(mag)	(mag)	(″)
2008/12/30.00	830.50	± 5	14.23 ± 0.01	$15\ 33\pm0\ 01$	15.07 ± 0.01	14.86 ± 0.01	14.86 ± 0.01	29
2008/12/50.00	833.48	$^{+9}_{+8}$	14.25 ± 0.01 14.39 ±0.01	15.55 ± 0.01 15.54 ± 0.01	15.07 ± 0.01 15.17 ± 0.01	14.65 ± 0.01	14.00 ± 0.01 14.98 ± 0.01	13
2009/01/01.98	840.52	±15	14.03 ± 0.01 15 47 ± 0.02	15.04 ± 0.01 15.76 ± 0.02	15.17 ± 0.01 15.21 ± 0.02	14.05 ± 0.01 14.80 ± 0.02	14.90 ± 0.01 14.87 ± 0.03	2.5
2009/01/03.02	855 45	+30	16.85 ± 0.02	16.76 ± 0.02 16.36 ±0.01	15.21 ± 0.02 15.38 ± 0.02	15.00 ± 0.02	14.87 ± 0.03 14.85 ± 0.01	$2.0 \\ 2.1$
2009/01/26.02	857 52	+30 + 32	16.00 ± 0.04 16.90 ± 0.04	16.00 ± 0.01 16.41 ± 0.02	15.30 ± 0.02 15.39 ±0.01	14.98 ± 0.01	14.85 ± 0.01 14.85±0.02	3.1
2009/01/27.01	858.33	+32	16.81 ± 0.02	16.41 ± 0.01	15.39 ± 0.01	15.00 ± 0.01	14.83 ± 0.01	2.8
2009/01/28.83	860.33	+35	16.83 ± 0.02	16.48 ± 0.01	15.00 ± 0.01 15.43 ± 0.01	15.00 ± 0.01 15.04 ± 0.01	14.89 ± 0.01	2.2
2009/01/30.88	862.38	+37	17.16 ± 0.04	16.53 ± 0.01	15.43 ± 0.01	15.02 ± 0.01 15.02 ± 0.01	14.85 ± 0.01	2.4
2009/02/02.89	864.39	+39	17.11 ± 0.04	16.53 ± 0.01	15.42 ± 0.01	14.99 ± 0.01	14.78 ± 0.01	2.0
2009/02/15.82	878.32	+53	17.58 ± 0.05	16.87 ± 0.01	15.43 ± 0.02	14.76 ± 0.01	14.82 ± 0.01	2.8
2009/02/16.96	879.46	+54	17.74 ± 0.06	16.84 ± 0.01	15.54 ± 0.01	15.09 ± 0.02	14.83 ± 0.01	3.1
2009/02/17.79	880.29	+55	17.79 ± 0.05	$16.87 {\pm} 0.01$	15.55 ± 0.01	15.10 ± 0.02	14.85 ± 0.01	3.1
2009/02/20.87	883.37	+58	$17.92 {\pm} 0.07$	$16.94{\pm}0.01$	$15.59 {\pm} 0.02$	15.09 ± 0.01	$14.84{\pm}0.01$	3.1
2009/02/23.77	886.27	+61	18.05 ± 0.04	$17.08 {\pm} 0.01$	_	15.13 ± 0.01	14.89 ± 0.01	3.1
2009'/02'/25.81	888.31	+63	$18.13 {\pm} 0.05$	$17.05 {\pm} 0.01$	$15.60 {\pm} 0.01$	15.14 ± 0.01	14.92 ± 0.01	2.9
2009/02/26.80	889.30	+64	_	_	15.62 ± 0.01	15.15 ± 0.01	_	3.8
2009/03/02.78	893.28	+68	_	$17.08 {\pm} 0.01$	$15.63 {\pm} 0.01$	$15.17 {\pm} 0.01$	$14.90 {\pm} 0.00$	4.4
2009/03/04.77	895.27	+70	_	17.15 ± 0.01	$15.66 {\pm} 0.01$	$15.18 {\pm} 0.01$	$14.94{\pm}0.01$	2.4
2009/03/09.76	900.27	+75	18.72 ± 0.14	$17.31 {\pm} 0.02$	$15.68 {\pm} 0.01$	_	$15.08 {\pm} 0.01$	2.3
2009/03/15.27	905.77	+80	_	$17.20 {\pm} 0.02$	_	15.21 ± 0.01	$14.92 {\pm} 0.01$	2.2
2009/03/24.80	915.30	+90	$18.73 {\pm} 0.08$	$17.42 {\pm} 0.01$	$15.94{\pm}0.01$	$15.48 {\pm} 0.01$	15.15 ± 0.01	2.5
2009/04/03.75	925.25	+100	19.28 ± 0.15	$17.66 {\pm} 0.02$	$16.10 {\pm} 0.01$	$15.57 {\pm} 0.01$	$15.31 {\pm} 0.01$	2.2
2009/04/05.61	927.22	+102	_	_	_	$15.67 {\pm} 0.01$	$15.41 {\pm} 0.01$	2.5
2009/04/10.67	932.33	+107	_	$18.38 {\pm} 0.06$	$16.76 {\pm} 0.01$	$16.20 {\pm} 0.01$	$15.84 {\pm} 0.02$	2.7
2009/04/17.60	939.25	+114	_	$19.83 {\pm} 0.06$	$17.91 {\pm} 0.01$	$17.13 {\pm} 0.01$	$16.68 {\pm} 0.01$	2.8
2009/04/22.68	944.11	+118	_	20.30 ± 0.16	$18.10 {\pm} 0.03$	$17.21 {\pm} 0.03$	$16.85 {\pm} 0.04$	3.5
2009/04/28.71	950.17	+125	_	$20.40 {\pm} 0.09$	_	$17.24 {\pm} 0.01$	$16.74 {\pm} 0.01$	2.2
2009/05/01.72	953.10	+127	_	20.22 ± 0.11	$18.06 {\pm} 0.04$	$17.27 {\pm} 0.04$	$16.70 {\pm} 0.04$	2.7
2009/05/23.69	975.18	+150	_	_	$18.30 {\pm} 0.04$	$17.47 {\pm} 0.02$	$17.00 {\pm} 0.03$	2.7
2009/05/25.70	977.21	+152	_	20.21 ± 0.12	$18.35 {\pm} 0.02$	$17.45 {\pm} 0.01$	_	2.6
2009/05/26.71	978.22	+153	_	20.15 ± 0.17	$18.33 {\pm} 0.04$	$17.49 {\pm} 0.02$	$17.06 {\pm} 0.03$	2.8
2009/06/14.69	997.19	+172	_	_	$18.39 {\pm} 0.04$	—	_	2.9
2010/01/17.98	1214.49	+389	_	_	_	$19.69 {\pm} 0.07$	$19.27 {\pm} 0.13$	1.9
2010/01/19.91	1216.42	+391	_	$21.47 {\pm} 0.12$	$20.34{\pm}0.05$	$19.76 {\pm} 0.05$	$19.52 {\pm} 0.07$	2.3
2010/02/13.91	1241.42	+416	_	-	$20.57 {\pm} 0.07$	$19.77 {\pm} 0.06$	-	2.2

NOTE. — The photometric observations are taken with the 1-m Sampurnan and Telescope, ARIES, Nainital. Errors in magnitude denote 1σ uncertainty.

^a With reference to the explosion epoch JD 2454825.6.

^c FWHM of the stellar PSF at V band. The nights for which we do not have any observation in V band, PSF is found through interpolation.

Long-slit low-resolution spectra (~ 6 to 14 Å) in the optical range $(0.33 - 1.0 \ \mu\text{m})$ were collected at eleven phases from +7d to +143d, including five phases from the 2-m IGO, three phases from the 9.2-m HET, two phases from the 6-m BTA and one epoch from the 3.6-m NTT ²². A journal of spectroscopic observations is given in Table 5. The spectroscopic data acquired from IGO, NTT and HET were reduced under the IRAF environment. Bias and flat-fielding were performed on all the frames. Cosmic-ray rejection was done using the Laplacian kernel detection method (van Dokkum 2001). All the data obtained from the BTA were reduced using programs in the IDL software environment.

The instrumental FWHM resolution of 2-m IGO spectra as measured from the [O I] λ 5577 emission skyline was found to lie between 6 and 10Å (~ 322 - 510 km s⁻¹). Flux calibration was done using standard spectrophotometric fluxes from Hamuy et al. (1994), assuming a mean extinction for the site. For HET, BTA and NTT, the resolution near 6000Å is about 10Å, 14Å, and 12Å respectively.

3. LIGHT CURVE EVOLUTION

According to theoretical interpretation, the entire broadband evolution of Type II SNe can be segmented into three phases: the rising phase, the plateau phase and the nebular phase.

3.1. The Rising Phase

The rising phase of the light curve is associated with the shock breakout phenomenon having been theoretically predicted for an SN explosion, in which, the radiation-dominated shock wave, generated through the reversal of iron core-collapse, starts to propagate outward through the onion shell-like layers of the progenitor; when the shock reaches regions with an optical depth of a few tens to unity, the radiation behind the shock escapes the outer surface giving rise to a hot $(T > 10^5 \text{ K})$ fireball which emits quasi-black body radiation in UV and soft-Xrays. This phenomenon is called shock breakout and depending on envelope mass, density structure and wind properties of the progenitor, the breakout light curve may last from a few hours to few a days (Grassberg et al. 1971; Chevalier 1976; Falk & Arnett 1977). Due to the short timescale, the detection of shock breakout is rare, and only recently were, the UV light curves of the entire

UT Date (yy/mm/dd)	JD 2454000+	Phase ^a (day)	J (mag)	$H \pmod{(\mathrm{mag})}$
2008/12/29.85	830.36	4.78	$14.66 {\pm} 0.054$	$14.55 {\pm} 0.064$
2009/01/03.87	835.37	9.78	$14.60 {\pm} 0.077$	$14.38 {\pm} 0.058$
2009/01/09.73	841.23	15.65	$14.39 {\pm} 0.059$	$14.16 {\pm} 0.072$
2009/01/14.73	846.23	20.65	$14.47 {\pm} 0.107$	$14.32 {\pm} 0.150$
2009/01/19.78	851.24	25.66	$14.54 {\pm} 0.150$	_
2009/01/24.76	856.26	30.67	$14.25 {\pm} 0.077$	$14.11 {\pm} 0.087$
2009/01/29.87	861.37	35.79	$14.37 {\pm} 0.050$	$14.14 {\pm} 0.065$
2009/02/05.66	868.16	42.58	$14.43 {\pm} 0.067$	$14.24 {\pm} 0.107$
2009/02/11.69	874.19	48.61	$14.20{\pm}0.107$	$14.59 {\pm} 0.241$
2009/02/16.72	879.22	53.64	$14.43 {\pm} 0.050$	$14.35 {\pm} 0.097$
2009/02/21.75	884.25	58.67	_	$14.18 {\pm} 0.063$
2009/02/27.75	890.25	64.67	_	$14.44 {\pm} 0.077$
2009/03/04.89	895.40	69.82	$14.40{\pm}0.074$	$14.12 {\pm} 0.085$
2009/03/10.74	901.24	75.66	_	$14.36 {\pm} 0.190$
2009/03/21.71	912.21	86.63	_	$14.50 {\pm} 0.090$
2009/03/26.73	917.23	91.65	_	$14.41 {\pm} 0.073$
2009/04/05.79	927.29	101.71	_	$15.00 {\pm} 0.147$
2009/04/10.80	932.30	106.71	_	$14.94{\pm}0.171$
2009/04/15.80	937.30	111.72	_	$14.86 {\pm} 0.173$
2009/04/20.52	942.02	116.44	_	$15.23 {\pm} 0.232$

TABLE 4 The REM near-infrared observation of SN 20081N.

^a With reference to the explosion epoch JD 2454825.6.

TABLE 5 Journal of spectroscopic observations of SN 20081N.

$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	UT Date (yy/mm/dd/hh.hh)	JD 2454000+	$\frac{\text{Phase}^{a}}{(\text{days})}$	$ \substack{ \text{Range} \\ \mu m } $	Telescope ^b	$ Grating (gr mm^{-1}) $	Slit width (")		Exposure (s)	${ m S/N^c}\ ({ m pix}^{-1})$
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	2008/12/31/22.80 2009/01/07/22.31 2009/02/17/19.74	832.45 839.43 880.32	+7 + 14 + 54	0.45 - 1.00 0.45 - 1.00 0.38 - 0.68	HET HET IGO	300 300 600	$1.0 \\ 1.0 \\ 1.5$	$5.0 \\ 5.0 \\ 1.4$	600 600 1800	$45 \\ 43 \\ 34$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	2009/02/22/22.50 2009/03/22/20.45	885.44 913.35	+60 +87	0.45 - 1.00 0.38 - 0.68	HET IGO	300 600	1.0 1.5	5.0 1.4	600 1800	35 38
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	2009/03/25/22.13 2009/04/03/19.14 2009/04/03/20.98	$916.42 \\ 925.26 \\ 925.33$	$^{+90}_{+99}_{+99}$	0.38 - 0.68 0.38 - 0.68 0.60 - 1.00	IGO IGO BTA	600 600 550	$1.5 \\ 1.5 \\ 2.1$	$1.4 \\ 1.4 \\ 3.5$	$2700 \\ 1800 \\ 3 \times 900$	35 35 60
2009/05/17/02.38 968.10 +143 0.33-0.80 NTT 300 1.0 4.0 2700 59	2009/04/22/15.49 2009/04/23/18.00 2009/05/17/02.38	$944.14 \\ 945.25 \\ 968.10$	$^{+118}_{+119}$ $^{+143}$	0.38 - 0.68 0.38 - 0.79 0.33 - 0.80	IGO BTA NTT	600 600 300	$1.5 \\ 2.1 \\ 1.0$	$1.4 \\ 3.5 \\ 4.0$	$1800 \\ 3 \times 900 \\ 2700$	$ \begin{array}{c} 11 \\ 65 \\ 59 \end{array} $

^a With reference to the explosion epoch JD 2454825.6.

^b HET : LRS on 9.2-m Hobby Eberly Telescope; IGO : IFOSC on 2-m IUCAA Girawali Observatory, India; BTA : SCORPIO on 6-m Big Telescope Alt-azimuth, Special Astrophysical Observatory, Russia; NTT : EFOSC2 on 3.6-m New Technology Telescope, ESO, Chile. ^c At 0.6 μ m.

shock breakout phase lasting several hours was reported for two Type IIP SNe – SNLS-04D2dc and SNLS-06D1jd (Gezari et al. 2008; Schawinski et al. 2008). Observation of the earliest UV and optical light curves of IIP SNe is crucial to model the shock breakout light curves and constrain the properties of SN progenitors (Tominaga et al. 2009). The earliest SNe IIP optical light curves have been studied in the past for SN 2005cs by Pastorello et al. (2009) and for SN 2006bp by Quimby et al. (2007) and the optical rising phase is rarely observed.

We model the early optical (ROTSE-IIIb R-band) light curves of SN 2008in using a simple model and put an observational constraint on the epoch and duration of shock breakout. Following the formulation of Waxman et al. (2007), it can be shown that just after the shock breakout, the intensity of the SN light at a fixed wavelength is proportional to the intensity of blackbody radiation at that wavelength (see Equation 1 of Cowen et al. 2010), while for the plateau phase Arnett (1980) derived an analytical expression (Equation 39 of that paper) and hence by combining these two equations one can approximate the overall intensity profile of a Type II SN during the rising phase through the plateau with the following expression :

$$I_{SN}(t) = \frac{A}{exp[B \times (t-t_0)^{0.5}] - 1} \times (t - t_0)^{1.6} + C \times exp[-\{(t/D) + (t/E)^2\}]$$
(1)

Here, the first term represents the phase associated with the shock breakout while the second term represents the plateau phase. Here 't' is the time measured in days since 2008 December 23.95 (JD = 2454823.5), the epoch when ROTSE had a non-detection of the SN with limiting magnitude 16.16 in the R band. The time of shock breakout, 't₀' along with the constants A, B, C, D and E are free parameters of the fit. The values of these constants depend on the nature of the progenitor and the kind of explosion. The above model



FIG. 3.— The ROTSE *R*-band light curve of SN 2008in. Reference time for x-axis is the epoch of ROTSE non detection on 2008 December 23.95. The evolution of post shock breakout flux has been modeled by using the simple analytical expressions given in Cowen et al. (2010) and Arnett (1980). The best fit time for shock breakout is marked with t_0 while the Δt_0 shows the duration of uncertainty. The t_d is the time of discovery reported by Nakano et al. 2008. The t_p is the time when *R* light shows the peak. The t_i is the time of inflection obtained using the procedure of Elmhamdi et al. (2003).

is fitted to the ROTSE data using the χ minimization technique (Press et al. 1992) and the fit is shown in Figure 3. We obtained a value of the shock breakout time as $t_0 = 2.13 \pm 0.83$ day. This value of t_0 is consistent with the first data point observed from ROTSE. We shall adopt t_0 , corresponding to JD = 2454825.6 as the explosion epoch for all the phases of SN 2008in. It is noted that at V and R bands, the SN was detected by Nakano et al. (2008), just 2.4 d after the shock breakout. We note, however, that depending on the extent of the envelope the true core reversal marking the SN explosion would have occurred a few hours earlier (e.g., SN 1987A) or a few days earlier for a red supergiant envelope ~ 50000 R_{\odot} (Quimby et al. 2007).

In Figure 3, ' t_p ' (+5.3 d) corresponds to the peak of the ROTSE *R*-band light curve. A similar peak in the UV light curves of the post-shock breakout phase of the Type IIP supernova SNLS-04D2dc lasting several days was identified as a secondary peak and this peak was explained as the shift of the spectral energy distribution toward longer wavelengths due to the rapid fall in temperature during free adiabatic expansion of optically thick plasma lead by a collisionless shock (Schawinski et al. 2008; Tominaga et al. 2009). For SN 2008in we identify the peak in the ROTSE R-band curve as a secondary peak and also speculate that for a few other type IIP SNe such as 1999em, 1999gi and 2005cs, where very tiny peaks before the plateau light curve were observed, such features are basically the footprints of the secondary peaks which could be a consequence of shock breakout.

3.2. The Plateau and Nebular Phase

The entire UV, optical and NIR light curves of SN 2008in are shown in Figure 4 and for comparison, the light curves of the archetypal low-luminosity SN 2005cs (Brown et al. 2009; Pastorello et al. 2009) are also plotted by scaling them in magnitude to match it with the observed plateau of SN 2008in. In the early plateau phase, a rapid drop in the UV flux and a slowly declining or constant optical and NIR flux are clearly apparent. The uvm2 magnitude declined from about 14.7 to 18.7 mag within 12 days. The decline rates of the SN 2008in flux in UVOT, U, B, and V bands are almost identical with that of SN 2005cs. Starting from the uvw2 until the R band, the measured decline rates from shock breakout to the plateau phase are approximately 0.32, 0.34, 0.25, 0.13, 0.04, 0.01 and $0.01 \text{ mag } d^{-1}$. The *IJH* light curves are almost flat even at early phases of the plateau. For



FIG. 4.— Light curves of SN 2008in in the UV, optical and NIR bands. The light curves of archetypal low-luminosity SN 2005cs are overplotted for comparison (data are taken from Brown et al. 2009; Pastorello et al. 2009). The light curves of SN 2005cs are scaled in magnitude to match the light curves of SN 2008in.

this SN, the plateau is well sampled in the BVRI bands, so we can accurately determine the plateau duration. As discussed in §3.1, the secondary peak ($t_p \sim +5.3$ d) in the prompt light curve is generated by the gradually cooling shock heated SN atmosphere and after that the plateau mechanism starts to dominate. On the other hand, between the plateau and the nebular phase there is another transitional state, when an inflection in the light curve can be seen (§7.1). Hence the plateau duration is precisely the time span between the secondary peak and the inflection. The inflection (t_i) in the V band light curve is observed to occur nearly 103.2 days after the shock breakout. Hence the time interval between the t_i and t_p is assumed as the duration of the plateau i.e. $(103.2 - 5.3) \sim 98$ days.

Once the hydrogen envelope is fully recombined and the ejecta becomes optically thin, the light curve enters into the nebular phase and it is sustained mainly by the energy output from the radioactive decays of the irongroup elements. During the plateau to nebular transition phase, the V-band magnitude drops from ~ 16.0 mag at around +90d (still in the plateau) to 18.1 mag at +122d (in the exponential light curve tail), i.e. ~ 2 mag in about one month. This drop is remarkably smaller than that of SN 2005cs, but consistent with the 2-3 mag drop observed in normal Type IIP SNe (Olivares E. et al. 2010). A linear fit to the tail from +120d to +400d gives the following decline rates [in mag (100d)⁻¹]: $\gamma_B \sim 0.44$, $\gamma_V \sim 0.84$, $\gamma_R \sim 0.91$, $\gamma_I \sim 1.03$ at B, V, R, and I bands which are similar to the values found in normal IIP SNe and comparable with the decay slope of ⁵⁶Co to ⁵⁶Fe, i.e. 0.98 mag (100d)⁻¹.

4. SPECTROSCOPIC EVOLUTION

The spectra of SN 2008in at 10 phases from +7d to +143d are shown in Figure 5. The spectra are corrected for the recessional velocity of the host galaxy ($\sim 1567 \pm 3$ km s⁻¹)²³. Spectral features are mainly identified as

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<sup>23</sup> http://leda.univ-lyon1.fr/
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per previously published line identifications for IIP SNe (Leonard et al. 2002; Pastorello et al. 2004). The two earliest spectra (+7d and +14d) show the blue continuum of blackbody emission and have broad P-Cygni profiles of H α , H β and He I λ 5876. The next two spectra (+54d and +60d) represent the mid-plateau phase and are marked by a decrease in the continuum and the appearance of more number of P-Cygni profiles for the permitted metallic (FeII, ScII, TiII, BaII, MgII), OI λ 7773, Na I D, and Ca II IR triplet lines similar to other normal Type IIP SNe (see Roy et al. 2011 and references therein). The spectra at phases +87d, +90d and +99drepresent the end stages of the plateau and these are marked by a redder continuum and decreasing line widths of hydrogen lines. The +118d, +119d and +143d spectra represent the nebular phase having negligible continuum and are marked by pronounced emissions of forbidden lines [CaII] $\lambda\lambda7291$, 7323, and fading of the absorption features of hydrogen and metallic lines.

In Figure 6, we compare the +99d spectrum of SN 2008 in with the spectra of other Type IIP SNe observed at similar phases that is at roughly the time when the light curve changes from the photospheric plateau to the nebular phase. The phases quoted for each event are with reference to the moment of inflection (ti; see Figure 3). The narrow P-Cygni H α profile and the presence of strong lines of BaII at $\lambda 6142$ and $\lambda 6497$ of SN 2008in show striking resemblance to the low-luminosity SNe 1997D, 1999br, 1999eu, 2001dc and 2005cs. On the other hand, the normal Type IIP SNe 2004et, 2008gz show broader profiles of $H\alpha$ and weaker lines of BaII. It is noted that in the blue wing of H α more metallic lines get resolved than that in the normal IIP SNe and this arises due to smaller line widths of hydrogen lines seen in low-luminosity IIP SNe (Fraser et al. 2010 and references therein). In addition, the Ba II line in the low-luminosity SNe show stronger absorption components due to lower ejecta temperature than that in the normal luminosity Type IIP SNe 2004et and 2008gz (Turatto et al. 1998). In Figure 7, we have identified the spectral features in a late-plateau phase (+99d) spectrum which covers the full wavelength range from 0.4 to 0.95 μ m. For line identification we have followed Pastorello et al. (2004), where different spectral lines were identified for the +102d spectrum of a low-luminosity SN 1999br. We are able to identify all the features and the spectral profiles of all the elements are similar to the archetypal low-luminosity IIP SNe 1997D, 2005cs and 1999br (Pastorello et al. 2004 and references therein).

Figure 8 shows the spectral evolution of H α , H β , NaID, FeII, BaII, and ScII lines. In Figure 9 (left panel), we show the expansion velocities of the ejecta derived from Balmer and FeII lines. The later is a good indicator of photospheric velocity. The expansion velocities of the H-envelope are estimated using absorption minima of the P-Cygni profiles and at the two earliest phases (+7d and +14d) H α and H β show a broad P-Cygni profile which, with time, becomes narrower at later phases keeping the position of the emission peak centered near zero velocity. The H α line velocity starts at about 7000 km s⁻¹ at +10d, reaches 4000 km s⁻¹ at +50d and flattens at a level of 1200 km s⁻¹ in the nebular phases. It can be seen from the right panel of Figure 9 that in the comparable phases, the H α line velocities of low-luminosity SNe are less than half those of normal IIP SNe 1999em and 2004et, whereas the SN 2008in velocities are more like low-luminosity SN 2005cs. To estimate the photospheric velocity of the transient we have computed the velocities of different Fe II lines λ 4924, λ 5018, and λ 5169 at different phases. The first marginal detection of these lines is in the +14d spectrum and they became prominent in later stages of evolution. For SN 2008in, the average velocity of these lines (and hence roughly the photospheric velocity) at +14d is about 4450 km s⁻¹, which is comparable with that of low luminosity SNe and much less than ordinary Type IIP events (see Figure 12 of Pastorello et al. 2009).

The emergence of absorption dips in the blue wings of $H\alpha$ is clearly seen in the post +60d spectra ²⁴ and the absorption dips have been marked with A, B and C in Figure 8. We distinguish 'A' as an absorption dip due to Ti II and H α and its equivalent width (EW) increases from 12.85 Å at +60d to 19.45 Å at +143d. This progressively stronger absorption dip requires a steep deceleration of the H α emitting zone along with its depletion by inner metal shells (like Ti II). The broader absorption dip 'B' is identified as a blend of BaII $\lambda 6497$, BaI, CaI and Ti II. An exceptionally strong Ba II lines is guite typical for cooler ejecta of less energetic fain SNe. The segment 'C' is speculated to be a footprint of Fe II multiplets as the evolution of its linewidths seems to be correlated with the other Fe II lines. It is noted, however, that in a few normal Type IIP SNe 1999em, 2004et and 2008gz, the spectral feature near 'C' at early plateau phases was also identified as a signature of high velocity $H\alpha$ component, and in the present spectra its presence cannot be completely ruled out.

The evolution of $H\beta$ is similar to the evolution of $H\alpha$. The P-Cygni profile of $H\beta$ is clearly visible in the +7d spectrum, and later on, the emission component is blanketed by various metal lines, mainly due to Fe II λ 4924, λ 5018 and λ 5169. Traces of H β and Fe II absorption dips indicate that from the beginning of the nebular phase H shells of ejecta and the regions containing Fe II ions move with a comparable velocity, reaching an asymptotic value $\sim 1000 \text{ km s}^{-1}$ nearly 140 days after the explosion. In the +14d spectrum, Fe II lines are marginally detected while in the spectra between +60d and +99d they are prominent. We investigate the temporal change of EW of a relatively less blended Fe II λ 4924 and from the spectra normalized to the H α peak value, we found that at an EW ~ 0.67 Å at +14d, increased to 17.45 Å at +54d and to 26.53 Å in the +99d spectrum. Later on, the EW decreased to 19.07 Å in the +118d spectrum and finally reaches to 18.27 Å in the +143d spectrum. This rapid fall in the EW of inner metal-rich shells plausibly indicates a decrease in the opacity of the SN inner ejecta.

Traces of the He I λ 5876 line are also seen in the +7d spectrum (Figure 8). The ratio of EWs between H α and He I for this spectrum is EW(H α)/EW(He I) ~ 1.99. In the +14d spectrum, this ratio has increased to 5.36. The steep decrement in EW of He I seems to be due to a quick recombination of He I ions as a result of the rapid fall

 $^{^{24}}$ This is the first time that the rapidly changing evolution of H α profiles during the late plateau to nebular phase has been so densely covered in a low-luminosity SN



FIG. 5.— Doppler corrected flux spectra of SN 2008in observed during +7 d to +143 d. Prominent lines are marked. Telluric features are marked with \oplus symbol. The sharp absorption dips seen near 5000 Å and 6500 Å in the +7d and +14d spectra are artifacts. The spectra signify the early photospheric phase (+7d, +14d), mid-plateau phase (+54d, +60d), late-plateau phase (+87d, +90d, and +99d) and nebular phase (+118d, +119d, and +143d).

in temperature of the constantly rarefying ejecta. From +54d, the emerging P-Cygni feature of NaI D becomes prominent. This feature seems to be a perfect P-Cygni throughout our entire spectral sequence. It simply indicates an uniform spherical distribution of NaI in the SN ejecta. In the three high S/N spectra labeled with +54d, +60d and +87d, a tiny absorption dip overlaid on the emission component of the NaI D P-Cygni is seen and since these spectra are redshift corrected, these dips are impressions of intervening interstellar matter present in the host galaxy. The absence of any similar absorption dip in the blue wing of the NaID profile, expected from the Galactic interstellar matter, confirms that there is little absorption from the Milky Way in the line of sight of SN 2008in.

The spectral evolution of Sc II λ 6248 and the s-process element Ba II λ 6142 is clearly seen in SN 2008in. In the +7d and +14d spectra, there is no clear evidence for

the presence of these two elements; however in the +54d spectrum they are prominent. The absorption components become stronger with time and persist until the +143d spectrum. In contrast, in the normal Type IIP SNe 2004et (Sahu et al. 2006), 1999em (Elmhamdi et al. 2003) and 2008gz (Roy et al. 2011) these features disappeared at ~ 170 day after the explosion. The ejecta of underluminous Type IIP SNe expand at a lower velocity than those of normal ones. So, the Ba lines in low luminosity events persist for a longer time just because the ejecta takes more time to cool-down. As a result it takes longer time to become optically thin.

In the nebular phase, the forbidden lines: [Ca II] $\lambda\lambda$ 7291, 7323, [O I] $\lambda\lambda$ 6300, 6364 and [Fe II] λ 7155 are among the strongest features visible in the spectra. Also, the line ratios of emission lines of Ca II and the O I doublet in late nebular phases is a good indicator of the progenitor mass (Fransson & Chevalier 1987). In Fig-



FIG. 6.— The end plateau spectrum (+99d) of SN 2008in, is compared with five low-luminosity SNe 1997D, 1999br, 1999eu, 2001dc, 2005cs (Pastorello et al. 2009 and references therein); two normal SNe 2004et (Misra et al. 2007), 2008gz (Roy et al. 2011) and a peculiar Type II SN 1998A (Pastorello et al. 2005), observed at comparable epochs. The phases quoted for each event are with reference to the moment of inflection (t_i ; see Figure 13), which marks the highest rate of decline in the V-band light curve during the end of the plateau and the beginning of nebular phase (Elmhamdi et al. 2003). All the spectra have low spectral resolution (~ 10 Å). The P-Cygni features of Na1 D, BaII λ 6142, ScII λ 6248, BaII λ 6497 and H α are marked.

ure 10, we compare the +139d and +283d spectra of SN 19997D (Benetti et al. 2001) with the nebular phase spectra of SN 2008in. The [Ca II] doublet is seen in the +99d, +119d and +143d spectra of SN 2008in, whereas a small footprints of the OI doublet can be seen in the +143d spectrum. Presence of [Fe II] line is not seen in our spectra. Considering that the explosion epoch of SN 1997D was quite uncertain, it is likely that the spectral evolution of H α and forbidden lines of SN 2008in is quite similar to that of SN 1997D.

5. EXTINCTION AND DISTANCE TOWARD SN 2008in

In order to determine the bolometric light curve and other physical parameters, a correct estimate of the extinction and distance toward the SN is essential. We adopt the Galactic reddening along the line-of-sight of SN 2008in as derived from the 100 μ m all sky dust ex-

tinction map (Schlegel et al. 1998), i.e. $E(B-V) = 0.0224 \pm 0.0003$ mag. In order to estimate reddening due to the host galaxy M61, we used the spectrum of SN 2008in taken on 2009 February 17 from the 2-m IGO telescope having good S/N (~ 40) and corrected for mean heliocentric radial velocity of the host ($cz \approx 1567 \pm 3$ km s⁻¹). Near the zero velocity, the spectrum showed a tiny absorption feature overlaid on the emission component of the P-Cygni profile of Na ID (§ 4). The EW of this absorption feature was computed as 0.535 ± 0.713 Å. The error in EW is estimated using Equation 6 of Vollmann & Eversberg (2006). It is known that the EW of the interstellar absorption bands is well correlated with the reddening E(B-V) estimated from the tail of SN Ia color curve (Barbon et al. 1990; Richmond et al. 1994; Turatto et al. 2003) and hence using the empirical rela-



FIG. 7.— The spectral line identifications in a late-plateau phase (+99d) spectra are shown. Spectral features are mainly identified as per previously published line identifications for IIP SNe (Leonard et al. 2002; Pastorello et al. 2004). Telluric features are marked with \oplus symbol.

tion E(B - V) = -0.01 + 0.16EW (EW in Å), given by Turatto et al. (2003), we obtain the host contribution as $E(B-V) \approx 0.076 \pm 0.104$, which is considerably higher than the Galactic contribution. This is consistent with the absence of any absorption feature due to Galactic interstellar matter in the blue wing of the emission component of Na ID profile. Consequently, we adopt E(B-V)(estimated as a sum of Galactic and host galaxy reddening) of 0.0984 ± 0.104 mag for SN 2008in. This corresponds to a visual extinction (A_V) of 0.305 ± 0.322 . Considering the uncertainty in the E(B-V) estimated using Na ID lines, a lower value of E(B - V) of 0.0448 ± 0.0006 mag (twice the Galactic reddening) and the corresponding A_V of 0.139 ± 0.002 cannot be ruled out. We will discuss its implication for the derived properties of the SN

The Hubble flow distance of the host galaxy M 61, after correction for the Virgo infall, is estimated as 13.7 ± 1.1

Mpc ²⁵. The distance estimated through the Tully-Fisher method is 12.1 ± 2.7 Mpc²⁶. Additionally, we also calculated the distance following the Standard Candle Method for Type II SNe (Hamuy & Pinto 2002; Hamuy 2005; Hendry et al. 2005). It is found that there is a strong correlation between the distance of Type IIP SNe along with their mid-plateau apparent V and I band magnitudes and the mid-plateau photospheric velocity, for a given cosmological model. For SN 2008in we estimate the mid-plateau (~ +50d) apparent magnitudes ~ 15.56 \pm 0.05 mag for V and ~ 14.77 \pm 0.06 mag for the I band. Considering the V and I band extinctions toward the SN, 0.305 ± 0.322 mag and 0.183 ± 0.193 mag respectively, we derive a distance of ~ 12.23 \pm 1.87 Mpc. The adopted value of mid-plateau photospheric velocity

²⁵ The cosmological model with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}, \Omega_m = 0.3$ and $\Omega_{\Lambda} = 0.7$ is assumed throughout the paper and the uncertainty corresponds to a local cosmic thermal velocity of 208 km s⁻¹ (Terry et al. 2002).

²⁶ http://nedwww.ipac.caltech.edu/



FIG. 8.— Spectral evolution of H β , FeII λ 4924, FeII λ 5018, HeI λ 5876, Na I D $\lambda\lambda$ 5890, 5896, Ba II λ 6148, Sc II λ 6248, Ba II λ 6497 and H α lines of SN 2008in during its transition from plateau to nebular phase. The dotted line at zero velocity corresponds to the rest wavelength, marked at the top of each panel. The sharp absorption dips seen in the H α profile for +7d and +14d spectra are artefacts.

is $2694.67 \pm 70 \text{ km s}^{-1}$ (see §7.2). Combining the above three measurements, we adopt the weighted mean distance of 13.19 ± 1.09 Mpc, which corresponds to a distance modulus of 30.6 ± 0.2 .

6. COLOR EVOLUTION AND BOLOMETRIC FLUX

Figure 11 shows the temporal evolution of the reddening-corrected broadband colors of SN 2008in. For comparison, we overplot the color curves for well studied SNe 1987A (Suntzeff & Bouchet 1990), 1999em (Elmhamdi et al. 2003), 2004et (Sahu et al. 2006; Maguire et al. 2010), 2005cs (Pastorello et al. 2009) and SN 2008gz (Roy et al. 2011). After small differences during the initial phases, the plateau phase color evolution of all Type II SNe is more or less similar. The $(U-B)_0$ and $(B-V)_0$ colors are blue during early photospheric phases and they become redder by about 1-2 mag in the plateau phase while the $(V-R)_0$ and $(V-I)_0$ colors evolves slowly and become red only by about 0.5 mag. The $(J-H)_0$ color remains constant at ~ 0.25

mag. During the transition phase from plateau to nebular, the low luminosity SN 2005cs showed a striking red peak in the $(B - V)_0$, $(V - R)_0$ and $(V - I)_0$ colors. For SN 2008in, this red peak is not present and its color evolution is found to be consistent with the normal Type IIP SNe.

The quasi-bolometric light curve estimated from the UV, optical and IR broadband (UVOIR) magnitudes of SN 2008in is shown in Figure 12 along with those of other Type II SNe. The extinction-corrected magnitudes are first converted into fluxes using zero-points given by Bessell et al. (1998) and then the total flux in UVOIR bands is obtained after a linear interpolation and integration between 0.203 and 1.67 μ m. The fluxes were calculated on those nights for which we had complete observations in UBVRI bands. For the initial two weeks the SN was detected by Swift/UVOT in near ultraviolet bands with a significant flux density. This contribution has been accounted for while measuring the net UVOIR flux. It has been assumed that contribution in the U band



FIG. 9.— Evolution of expansion velocity of SN 2008in. The left panel shows the velocity of different elements in the ejecta of SN 2008in. The black filled circle shows the estimated value of mid-plateau photospheric velocity, calculated using the relation mentioned in Roy et al. (2011). The right panel shows a comparison of velocity derived from H α with that of low-luminosity event like SN 2005cs (Pastorello et al. 2009) and normal Type IIP SNe 1999em (Elmhamdi et al. 2003) and 2004et (Sahu et al. 2006).

is mainly important during the plateau and decreases rapidly to about 5% in the nebular phase (see Misra et al. 2007; Roy et al. 2011 and references therein). From the light curve it is noticeable that the near ultraviolet flux is almost negligible beyond 20 days after explosion. The JH contribution during the plateau is calculated from our data, whereas for the nebular phase, due to lack of data, we are not able to make any direct measurement. For most low luminosity SNe IIP, the NIR flux contribution in the nebular phase is about 50% of the total flux (Pastorello et al. 2009). We have therefore increased the net flux by 55% to account for the maximal contributions from the U and NIR bands at phases later than +140d. Figure 6 and 12 clearly demonstrate that spectroscopically SN 2008in appears to be like low-luminosity SNe IIP, but photometrically it appears quite normal. In the following section, the UVOIR curve is used to estimate the amount of radioactive ⁵⁶Ni and other physical parameters that characterize the explosion and the progenitor star.

7. PHYSICAL PARAMETERS

The physical entities that seem to govern the entire scenario are mainly associated with the nature of the progenitor and the radioactive elements (mainly ⁵⁶Ni), generated inside the inner portion of the ejecta during the explosion. ⁵⁶Ni is synthesized by the explosive burning of Si and O during the shock breakout (Arnett 1980, 1996). Over time this material is eventually converted to ⁵⁶Co and then to ⁵⁶Fe by means of radioactive transitions having decay times of 8.77 and 111.3 days respectively. The γ -rays and neutrinos emitted during this process sustain the nebular phase light curve and consequently, the observed tail luminosity becomes a good tracer for the ejected synthesized ⁵⁶Ni.

7.1. Produced Radioactive Nickel

We use different methods to estimate mass of 56 Ni. Hamuy (2003) proposed a relation between bolometric tail luminosity and the synthesized ⁵⁶Ni during core collapse SNe, by considering the underlying assumption that all γ -rays emitted during the radioactive decay make the ejecta thermalised. For SN 2008in, the time interval spanned by the observations is about 416 days, where the first ~ 100 days are reserved for photospheric evolutionary processes. The average V band magnitude during the nebular phase calculated using the data between +114d and +416d is \sim 18.91, which corresponds to the V magnitude at $\sim +222$ d. Taking the extinction correction as $(A_V = 0.305 \pm 0.322 \text{ mag}; \S5)$, a bolometric correction of 0.26 ± 0.06 mag (Hamuy 2001) and a distance modulus 30.6 ± 0.2 , the derived tail luminosity at this fiducial time is $(3.02 \pm 1.95) \times 10^{40} \,\mathrm{erg \, s^{-1}}$. Within the errors, this value is consistent with the bolometric flux at a comparable epoch, determined in $\S6$. This implies that the amount of ⁵⁶Ni produced is this process is $M_{\rm Ni} = 0.0157 \pm 0.0102 M_{\odot}.$

All Type IIP SNe show an inflection in the light curve during the transition from the plateau to the nebular phase. Statistically it has been shown that the steepness of the V-band light curve slope (defined as $S=dm_V/dt$) at the inflection time (t_i) is anti-correlated with ⁵⁶Ni mass (Elmhamdi et al. 2003). For SN 2008in, the V band light curve with its well-sampled transition phase, shows a value of steepness $S = 0.151 \pm 0.044$ mag d⁻¹ (Figure 13) and the epoch of inflection is $t_i \approx \pm 101.5d$ with respect to date of discovery. This corresponds to $M_{\rm Ni} = 0.0175 \pm 0.002 M_{\odot}$. This result is consistent with the value measured using the Hamuy (2003) scheme. According to Elmhamdi et al. (2003), the amount of ⁵⁶Ni



FIG. 10.— Forbidden lines during the nebular phases of SN 2008in are compared with the low-luminosity SN 1997D. The lines [OI], [FeII] and [CaII] are marked.

can also be derived from the plateau absolute V-band magnitude using the relation $\log M_{\rm Ni} = -0.438 M_V(t_i - 35) - 8.46$. Here $M_V(t_i - 35)$ is the absolute V magnitude 35 days prior to the day of inflection. For SN 2008in, $M_V(t_i - 35) \approx -15.23$, which again corresponds to a ⁵⁶Ni mass around $0.016 M_{\odot}$.

Comparison of the tail luminosity with that of SN 1987A – a well studied proximate event, is also used for estimation of the ⁵⁶Ni mass. A linear least square fit on the nebular light curve tail shows that at +222d the luminosity of SN 2008in is about $1.54 \times 10^{40} \,\mathrm{erg \, s^{-1}}$, while SN 1987A had a luminosity ~ $1.10 \times 10^{41} \,\mathrm{erg \, s^{-1}}$ (Figure 12). Since the ⁵⁶Ni mass produced by SN 1987A is about $0.075 M_{\odot}$, the amount of ⁵⁶Ni in the case of SN 2008in is [$(1.54/1.10) \times 10^{-1}$] × $0.075 \approx 0.0105 M_{\odot}$.

The above estimates are consistent with each other and hence we adopt a mean value for 56 Ni mass as $0.015 \pm 0.003 M_{\odot}$

7.2. Explosion Energy and Mass of Progenitor Star

We use the radiation-hydrodynamical simulations of core-collapse IIP SNe by Dessart et al. (2010) to infer the explosion energy $(E_0$ - kinetic plus thermal, expressed hereafter in units of 10^{51} erg or foe). These simulations suggest that in a given progenitor larger explosion energies yield larger ejecta velocities and this implies that the ejecta kinematics can be used to put constraint on E_0 . For SN 2008in, the expansion rate of the H-rich progenitor envelope as derived from absorption minima in $H\alpha$ at 15d after shock breakout is about 6300 km s⁻¹ (see Figure 8), which is higher than 4700 km s⁻¹ (for SN 2005cs; Pastorello et al. 2009) and lower than 8800 $\rm km \, s^{-1}$ (for SN 1999em; Elmhamdi et al. 2003). For the SNe 2005cs and 1999em, the simulation results (for explosion of a non-rotating solar metallicity pre-SN stars) predicted $E_0 \sim 0.3$ foe and $\gtrsim 1$ foe respectively and these values are found to be consistent with that determined from the actual hydrodynamical modeling of the SN light curves, e.g. 0.4 foe (Utrobin & Chugai 2008) and 0.3 foe (Pastorello et al. 2009) for SN 2005cs; 1.3 foe (Utrobin



FIG. 11.— Colour curve of SN 2008in is shown with filled circles. Also shown are the SNe 1987A (solid line), 1999em (open triangle), 2004et (filled triangle), 2005cs (open circle) and 2008gz (open square).

2007) and 1.25 foe (Bersten et al. 2011) for SN 1999em. For SN 2008in, the simulations suggest $E_0 \sim 0.5$ foe.

Accurate determinations of explosion parameters such as E_0 , the ejected mass $(M_{\rm ej})$ and the pre-SN radius (R_0) of the progenitor require detailed hydrodynamical modeling of the light curves and spectra and this being a non-trivial task (beyond the scope of this paper) has only been attempted for a few SNe. In order to have an estimate of explosion parameters of SN 2008in here, we employ the analytical relations derived by Litvinova & Nadezhin (1985) correlating the observed parameters $(M_{Vmp}$ - the mid-plateau absolute magnitude at V, v_{mp} - the mid-plateau photospheric velocity and Δt_p -the plateau duration) with the physical parameters $(E_0, M_{ej} \text{ and } R_0)$ based on a grid of hydrodynamical models for different values of physical parameter for Type IIP SNe. We note, however, that these approximate formulae have limitations owing to the poorlymeasured observables and the simplified physical conditions such as non-inclusion of the effect of nickel heating, use of old opacity tables, neglecting the effect of line opacity and using outdated pre-SN models (Smartt et al. 2009; Bersten et al. 2011 and references therein). As a result, only approximate values of the physical parameters can be derived using these relations. Fortunately, the observed parameters are derived very accurately for SN 2008in. The Δt_p is ~ 98 days (§3.2), the $v_{\rm mp}$ is 2694 ± 70^{27} km s⁻¹ and the M_{Vmp} is estimated as -15.32 ± 0.38 mag²⁸ Now, employing analytical relations, we estimate $E_0 \sim 0.54$ foe, $M_{\rm ej} \sim 16.7 M_{\odot}$ and $R_0 \sim 127 R_{\odot}$. The explosion energy derived in this way is consistent with that predicted from the hydrodynamical

 27 This is the mean value of the velocities computed from the lines of Fe II $\lambda4924,\,\lambda5018$ and $\lambda5169$ in the +54d spectrum.

 28 These observed values of M_{Vmp} and v_{mp} can be compared with the estimates derived empirically between the mass of $^{56}{\rm Ni}$, and the v_{mp} and M_{Vmp} for a larger sample of IIP SNe (see Section 8.3 of Roy et al. 2011). Using the $^{56}{\rm Ni}$ mass of $0.015\pm0.003M_{\odot}$ we find $v_{mp}=2916\pm220\,{\rm km\,s^{-1}}$ and $M_{Vmp}=-15.37\pm0.23$ mag which are consistent with those measured observationally.



FIG. 12.— Comparison of the quasi-bolometric light curve of SN 2008in with the low-luminosity SNe 1997D, 2005cs; the normal SNe 1999em, 2004et, 2008gz; and the peculiar Type II SN 1987A.

simulations of the ejecta kinematics.

The explosion energy of SN 2008in indicates that the event was less energetic than the standard IIP SNe 1999em, 2004et and more energetic than SN 2005cs. Assuming a net mass loss ~ $0.5M_{\odot}$ due to stellar wind ²⁹ and accounting for a compact remnant with mass ~ $1.5 - 2M_{\odot}$ (Sahu et al. 2006), we find that the initial mass of the progenitor to be $\leq 20M_{\odot}$.

8. DISCUSSION IN CONTEXT OF OTHER TYPE IIP SNE

In Figure 14, we compare the light curve of SN 2008in in absolute V-band magnitude with a sample of 13 other well studied Type II SNe taken from the literature (see Misra et al. 2007 and Roy et al. 2011 for the references). The sample includes two low-luminosity Type IIP SNe 1997D, 2005cs; nine normal Type IIP SNe 1990E, 1992H, 1999gi, 1999em, 2003gd, 2004A, 2004dj, 2004et, 2008gz and two peculiar Type II SNe 1987A, 1998A. It is seen that the Type II SNe show a wide range of mid-plateau luminosity, i.e. from -14 to -17 mag. The data on low luminosity SNe (see also Table 4 in Pastorello et al. 2006) indicate that their M_{Vmp} lie between -14 and -15 mag, with the only exception of SN 2002gd, whereas the other SNe are seen to lie between -16 and -17 mag. With the M_{Vmp} of -15.32 mag³⁰, the SN 2008in present another case which occupies the gap between a low and normal luminosity Type IIP events.

The tail luminosity is an important indication of ejected nickel mass. The average tail luminosity of SN 2008in (Mv ~ -12.16) is nearly 2.3 mag brighter than those of the low luminosity SNe 1997D and 2005cs (~ -9.6 mag), and about 1.3 mag fainter than the nor-

 $^{^{29}}$ SN 2008 in was not detected in X-ray. The lack of X-ray emission can be used to constrain the mass-loss rate of the progenitor system that could be heated by the outgoing shock to X-ray emitting temperatures. Following the discussion in Immler et al. (2007) and references therein, an upper limit to the pre-SN mass-loss rate of $5 \times 10^{-6} \ M_{\odot} \ {\rm yr^{-1}} \ (v_{\rm w}/10 \ {\rm km \ s^{-1}})$ with an uncertainty of a factor of two to three is obtained.

 $^{^{30}}$ Considering a lower value of $A_V=0.139\pm 0.002,$ we get $M_{Vmp}=-15.1\pm 0.2$



FIG. 13.— Determination of steepness parameter from the apparent V-band light curve of SN 2008in. See text for details.

mal Type IIP SNe 1992H, 2004et and 2008gz (~ -13.5 mag). A close inspection shows that the average tail luminosity of SN 2008in is roughly similar to Type IIP SNe 2004dj in NGC 2403 and 2003gd in M74, whereas their plateau is brighter (~ 0.5-0.8 mag at mid-plateau) than SN 2008in. These comparison of nebular phase luminosities with other SNe are consistent with a quite modest radioactive ⁵⁶Ni production for SN 2008in (§7.1).

The luminosity and shape of the tail depend on the ⁵⁶Ni mass and the radiant energy per unit ejected The first parameter determines the mass $(E/M_{\rm ei})$. absolute magnitude, while the second determines its decay rate (see Turatto et al. 1998 and references therein). The measured value of ejected ⁵⁶Ni mass for SNe 2003gd, 2004dj and 2008in are respectively ~ 0.016 (Hendry et al. 2005), 0.017 (Vinkó et al. 2006) and $0.015 M_{\odot}$. These are more than twice the ⁵⁶Ni amount synthesized by the low-luminosity SNe 1997D $(\sim 0.002 M_{\odot})$, Turatto et al. 1998) and 2005cs (0.003 - $0.004 M_{\odot}$, Pastorello et al. 2009), although at least three times less than the 56 Ni produced by normal Type IIP SNe such as 1992H ($\sim 0.075 M_{\odot}$, Clocchiatti et al. 1996), 2004et (~ $0.06M_{\odot}$, Sahu et al. 2006; Maguire et al. 2010) and 2008gz (~ $0.05M_{\odot}$, Roy et al. 2011). Similarly, the measured value of $E/M_{\rm ej}$ ratio for SN 2008in is larger than that for the low-luminosity SN 1997D and

lower than other Type IIP events 31 .

Though we have an approximate value of the pre-SN radius (128 R_{\odot}), it can be seen that the progenitor radius of SN 2008in is smaller than or comparable to that of low-luminosity SNe 1997D (~ $300R_{\odot}$, Turatto et al. 1998), 2005cs (100 - $600R_{\odot}$, Pastorello et al. 2009, Utrobin & Chugai 2008), 1999br (114 R_{\odot} , Zampieri et al. 2003) and much smaller than that of normal Type IIP SNe like 1992H (> $600R_{\odot}$, Clocchiatti et al. 1996), 2004et (~ $530R_{\odot}$, Misra et al. 2007). We also note that SN 2008in progenitor is only slightly larger than the blue supergiant progenitors of SNe 1987A and 1998A (Pastorello et al. 2005).

Information about the metallicity of the SNe location is essential to constrain the triggering mechanism of the SN explosion (Heger et al. 2003). The oxygen abundance [O/H] of SN 2008in is about 8.44 dex ³² which is marginally lower than the solar value of 8.65 dex (Asplund et al. 2009). A comparison ³³ of [O/H] of all the events in our sample indicate that there is no clear trend in the type of event and the metallicity.

The photometric and spectroscopic comparisons of SN 2008in with some Type II SNe covering a wide range of physical parameters puts observational constraints on the nature of the progenitor of this event, pointing toward a star that was more compact than a typical M-Type red supergiants, and closer to a blue supergiant, may be an yellow supergiant. However, the direct detection of the progenitors of a few faint SNe IIP in preexplosion archive Hubble Space Telescope images seems to contradict this conclusion (see Smartt et al. 2009 for a review). Our spectroscopic study suggests that in SN 2008in the hydrogen envelope ejected by the explosion is smaller than in most Type II events but larger than that ejected in low luminosity events (such as SN 1997D, SN 1999br and SN 1999eu). This may be due to a significant mass loss of the parent star in the latest stages of its life. This conclusion is also partially supported by the upper limit mass loss rate of the progenitor revealed from the X-ray study $(\S7.2)$. In view of the upper mass limit of $20 \ M_{\odot}$ progenitor, occurrence of the event in a sub-solar metallicity region and following the evolution of a sin-

³¹ The measured value of the E/M ratio for SN 2008in is 3.5×10^{49} erg M_{\odot}^{-1} , which is twice the E/M ratio measured for SNe 1997D (1.7×10^{49} erg M_{\odot}^{-1} , Turatto et al. 1998), but it is half those of SN 2008gz (7.3×10^{49} erg M_{\odot}^{-1} , Roy et al. 2011), SN 2004dj (7.3×10^{49} erg M_{\odot}^{-1} , Vinkó et al. 2006) and SN 2004et (7.5×10^{49} erg M_{\odot}^{-1} , Sahu et al. 2006). Finally, it is barely a third the E/M ratios of SNe 1987A (11.3×10^{49} erg M_{\odot}^{-1} , Hamuy 2003) and 2003gd (13.0×10^{49} erg M_{\odot}^{-1} , Hendry et al. 2005), and about seven times smaller than the luminous event like SN 1998A (25.4×10^{49} erg M_{\odot}^{-1} , Pastorello et al. 2005). ³² Derived from the [O/H](= $12 + log(N_O/N_H)$)- M_B relation given by Pilyugin et al. (2004) for a given deprojected radius and

³² Derived from the $[O/H](= 12 + log(N_O/N_H))$ - M_B relation given by Pilyugin et al. (2004) for a given deprojected radius and galaxy type. The host of SN 2008in is a spiral galaxy of SBbc Type and the SN location is ~ 1.8' away from the center of the host. This corresponds to a deprojected distance ~ 7.03 kpc.

 33 We have taken the [O/H] value from Smartt et al. (2009), though for a few cases the values are estimated using the [O/H]- M_B relation mentioned in Pilyugin et al. (2004). The [O/H] for normal Type IIP SNe 2008gz, 1999em and low-luminosity SNe 1999eu, 2005cs is about 8.6. For the low-luminosity SNe 1999k, 1999br, 2001dc, 2003gd and 2004dj, 1997D and the normal SN 2004et, it is about 8.4. In many cases, we have adopted the [O/H] value from Smartt et al. (2009). For some cases we have estimated the value by adopting the [O/H]- M_B relation mentioned in Pilyugin et al. (2004).



FIG. 14.— A comparison of the absolute V-band light curve of SN 2008in with the low-luminosity SNe 1997D, 2005cs; the normal SNe 1999em, 2004et, 2008gz, 2004dj, 2004A, 2003gd, 19992H, 1999E, 1999gi; and the peculiar Type II SNe 1987A, 1998A. The decline rate of emission exptected from radioactive decay of 56 Ni to 56 Co to 56 Fe is shown with dot-dashed line.

gle massive star as a function of metallicity (Heger et al. 2003), we rule out the possibility of "fall back of ejecta to BH" scenario in case of SN 2008in, supporting the scenario of a weak explosion of a relatively compact progenitor.

9. CONCLUSION

Low-luminosity Type IIP SNe belong to a poorly known class of events due to the unknown nature of their progenitors as well as the explosion mechanism. Spectroscopic as well as photometric characteristics of these events are significantly different from the normal and peculiar Type IIP SNe. The number of such underluminous events discovered so far is relatively small. This could plausibly be a selection effect. If we confine the search of core-collapse SNe to a small volume, the majority of SNe discovered in the local universe are of Type IIP ($\sim 48.2\%$, Smith et al. 2011) and probably a large fraction of them would turn out to be underluminous. Incidentally SN 1999br, SN 2005cs and SN 2009md are the only events of this group to have been discovered soon after the corecollapse and whose data are publicly available.

In this study we have reported the results of an extensive photometric and spectroscopic follow-up campaign for SN 2008in in the ultraviolet, optical and near-infrared domains. The SN was also observed (although not detected) at X-ray and radio wavelengths. The SN was observed soon after its explosion, during the fast rise of the light curve to the optical maximum. Evidence for a shock breakout in SN 2008in was primarily derived through the analysis of the ROTSE-IIIb R-band light curve. An upper limit of 16.16 mag in the R-band was estimated about two days before the discovery.

Modeling the R-band light curve allowed us to estimate a reliable epoch of the shock breakout, with an uncertainty of about two days. The plateau phase in SN 2008in lasts about 98 days which is marginally shorter than in normal Type IIP SNe. The evolution of the bolometric light curve indicate that the event is somewhat in between the normal and faint Type IIP event.

The luminosity of the nebular phase light curve indicates an ejected 56 Ni mass of $\sim 0.015 M_{\odot}$, a factor two higher than that derived for low-luminosity IIP SNe. The spectroscopic evolution of SN 2008in is similar to those of low-luminosity IIP SNe (1997D, 1999br and 1999eu), indicated particularly by the strong presence of Ba II lines, the narrow line widths of $H\alpha$ lines, and the expansion velocities of SN during the photospheric phases (~ 3000 km s⁻¹) and the nebular phases ($\sim 1200 \text{ km s}^{-1}$). The ejecta kinematics of SN 2008in are consistent with lessenergetic (~ 5×10^{50} erg) Type IIP SNe.

Spectroscopically SN 2008in appears to be like lowluminosity SNe IIP, but photometrically it appears close to a normal type IIP event.

From the light curve and spectra of SN 2008in, we could determine quite accurate values of the observed properties such as plateau duration, mid-plateau luminosity as well as the photospheric velocity and this has helped us to comment on the properties of the explosion and the progenitor star. Using semi-analytical formulae by Litvinova & Nadezhin (1985), we could estimate approximate values of the explosion energy $\sim 5.4 \times 10^{50}$ erg, the ejected mass $\sim 17 M_{\odot}$ and the pre-SN radius ~ $127R_{\odot}$. The explosion energy of SN 2008in is smaller than the normal ($\geq 10^{51}$ erg) Type IIP events, although higher than that estimated in underenergetic (~ 10^{50} erg) SNe IIP. We could provide an upper limit to the mass loss rate of the progenitor as $5 \times 10^{-6} M_{\odot} \text{ yr}^{-1} (v_{w}/10 \text{ km s}^{-1})$ whereas the upper limit for the main-sequence mass of the progenitor star is estimated as 20 M_{\odot} .

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