# A method to determine distances to molecular clouds using near-IR photometry 

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#### Abstract

Aims. We aim to develop a method to determine distances to molecular clouds using $J H K$ near-infrared photometry. Methods. The method is based on a technique that aids spectral classification of stars lying towards the fields containing the clouds into main sequence and giants. In this technique, the observed $(J-H)$ and $\left(H-K_{\mathrm{s}}\right)$ colours are dereddened simultaneously using trial values of $A_{V}$ and a normal interstellar extinction law. The best fit of the dereddened colours to the intrinsic colours giving a minimum value of $\chi^{2}$ then yields the corresponding spectral type and $A_{V}$ for the star. The main sequence stars, thus classified, are then utilized in an $A_{V}$ versus distance plot to bracket the cloud distances. Results. We applied the method to four clouds, L1517, Chamaeleon I, Lupus 3 and NGC 7023 and estimated their distances as $167 \pm 30,151 \pm 28,157 \pm 29$ and $408 \pm 76 \mathrm{pc}$ respectively, which are in good agreement with the previous distance estimations available in the literature.


Key words. dust, extinction - ISM: clouds - stars: distances - infrared: ISM

## 1. Introduction

The knowledge of the distance to interstellar clouds which are either isolated or associated with large complexes, star forming or starless, is important to provide observational constraints to the properties of clouds and pre-main sequence (PMS) stars (e.g., Yun \& Clemens 1990; Clemens et al. 1991; Kauffmann et al. 2008). However, distances to many of these clouds, especially those that are isolated, are highly uncertain (in some cases, by a factor of 2, see Hilton \& Lahulla 1995).

In recent years, successful efforts have been made to estimate distances to large complexes using a number of methods. For example, combining extinction maps from the Two Micron All Sky Survey (2MASS) with Hipparcos and Tycho parallaxes, Lombardi et al. (2008) estimated accurate distances to the Ophiuchus and Lupus dark cloud complexes; the Hipparcos parallax measurements and B-band polarimetry were utilized by Alves \& Franco (2007) to estimate the distance of the Pipe nebula. Using multi-epoch very long baseline array (VLBA) observations of a number of young stellar objects (YSOs), Loinard et al. (2007) measured the trigonometric parallaxes and hence distances with a precision of about $1-4 \%$ (Loinard et al. 2007, 2008). The distances to the parent clouds that harbour these YSOs are also constrained with similar precision.

However, it is difficult to establish distances to isolated clouds devoid of low-mass YSOs and primary indicators such as ionizing stars or reflection nebulae. Distances to these clouds can be assigned through their association (in position-velocity space) with nearby larger molecular clouds (e.g., Kauffmann et al. 2008) which have their distances determined through the
abovementioned methods. The traditional ways of determining distances to isolated clouds utilize the star count method (Bok \& Bok 1941) or Wolf diagrams (Wolf 1923), which is limited due to its dependence on the questionable extrapolation of luminosity functions. Additional methods of assigning distances to small dark clouds have involved bracketing the cloud distance by using spectroscopic distances to stars (Hobbs et al. 1986) and using the equivalent widths of interstellar Ca II H and K lines of stars close in front of and behind the cloud (Megier et al. 2005). However, spectroscopy of a large number of stars projected on a cloud is extremely demanding in terms of telescope time. Because of the absence of nearby associated large molecular clouds or stellar spectroscopic data, distances of most of the isolated small clouds remain unknown.

Distances to clouds can be determined by bracketing them using spectral classification of stars projected on to the cloud purely by photometry carried out in the Vilnius photometric system (Straižys 1991; Straižys et al. 1992). The distance and reddening for individual stars can also be obtained using the very well calibrated Stromgren uvby - $\beta$ intermediate-narrow- band photometric system. The physical parameters of stars thus obtained are utilized to bracket the cloud distances (e.g., Nielsen et al. 2000; Franco 2002). Another innovative method was developed by Peterson \& Clemens (1998) by identifying M dwarfs lying both in front of and behind the cloud. The reddening and spectral types of these M dwarfs were determined from photometry alone due to their conspicuous position in the $(V-I)$ versus ( $B-V$ ) colour-colour (CC) diagram. The optical $V, R, I$ and near-IR (NIR) $J, H$ and $K$ 2MASS colours were utilized by Maheswar et al. $(2004,2006)$ to bracket cloud distances, again
by spectral classification of stars projected on to them from photometry alone.

The enormous NIR dataset provided by the 2MASS (Kleinmann et al. 1994) was extensively used by several authors, mainly (a) to discriminate between normal stars and stars with significant NIR emission from circumstellar material (e.g., Itoh et al. 1996) and (b) to map the extinction in dense molecular clouds (Cambrésy et al. 2002; Lombardi et al. 2006). Dutra et al. (2003) built an $A_{K}$ extinction map of a $10^{\circ}$ field towards the Galactic centre using 2MASS $J$ and $K_{\mathrm{s}}$ magnitudes. They utilized the upper giant branch of colour-magnitude diagrams and their dereddened mean locus built from previous studies of bulge fields to estimate the extinction. In this work, we present a method by which distances to molecular clouds are estimated using $(J-H),\left(H-K_{\mathrm{s}}\right)$ colour indices of the main sequence stars in the spectral range from A0 to K7. The technique is based on a methodology that provides spectral classification of stars projected on to the fields containing the cloud into main sequence and giants by simultaneously dereddening the $(J-H)$ and ( $H-K_{\mathrm{s}}$ ) observed colours using trial values of $A_{V}$ and a normal interstellar extinction law. The main sequence stars are then utilized in an $A_{V}$ vs. distance (d) plot to bracket the cloud distances. In Sect. 2 we describe the method for finding the spectral type, $A_{V}$, and distances to the molecular clouds using NIR data from the 2MASS. The uncertainties involved in the determination of the $A_{V}$ values and the distances from the method are also discussed in this section. In Sect. 3 we describe the criteria used to extract the data from the 2MASS database. In Sect. 4 we apply our technique to four clouds, L1517, Chamaeleon I, Lupus and NGC 7023. We discuss the method in determining distances to dark molecular clouds in Sect. 5. Finally, we conclude the paper by summarizing the results obtained in Sect. 6.

## 2. The method

The photometric distance $d$ to a star is estimated using the distance equation

$$
\begin{equation*}
d(p c)=10^{\left(V-M_{V}+5-A_{V}\right) / 5}=10^{\left(K-M_{K}+5-A_{K}\right) / 5} \tag{1}
\end{equation*}
$$

where $V, M_{V}, A_{V}, K, M_{K}$, and $A_{K}$ are the apparent magnitude, absolute magnitude and extinction in $V$ and $K$ filters respectively. Spectral type and luminosity class information are needed to determine both absolute magnitude of a star and the extinction suffered by it due to interstellar dust in the line of sight. Spectral type and luminosity class also determine the various colours of a star. Therefore, in principle, one could determine the spectral type and the luminosity of a star from its observed colours provided the extinction is zero. In practice, observed colours are reddened due to interstellar extinction, which is wavelengthdependent. By assuming a value for $A_{V}$ and an extinction law one could estimate the colour excesses and correct the observed colours to find the intrinsic colours of the stars and hence their spectral types.

### 2.1. The $(J-H),\left(H-K_{s}\right)$ intrinsic colours

In Fig. 1, we present the $(J-H),\left(H-K_{\mathrm{s}}\right) \mathrm{CC}$ diagram. The solid line represents locations of unreddened main sequence stars. Usually the unreddened colours of different spectral types and luminosity classes in the 2MASS system are obtained from Koornneef (1983) or Bessell \& Brett (1988) using the transformation equations given by Carpenter (2001). We produced the intrinsic colours of the main sequence stars in the spectral range


Fig. 1. The $(J-H)$ vs. $\left(H-K_{\mathrm{s}}\right)$ CC diagram drawn for stars towards a region chosen arbitrarily on the sky to illustrate the method described in Sect. 2. The solid curve represents locations of unreddened main sequence stars. The reddening vector for an A0V type star drawn parallel to the Rieke \& Lebofsky (1985) interstellar reddening vector is shown by the dashed line. The locations of the main sequence stars of different spectral types are marked with square symbols. The region to the right of the reddening vector is known as the NIR excess region and corresponds to the location of PMS sources. The dash-dot-dash line represents the loci of unreddened CTTSs (Meyer et al. 1997). The closed circles represent the observed colours and the arrows are drawn from the observed to the final colours obtained by the method for each star.
from A0 to K7 in the 2MASS system directly from the observations by selecting stars from the All-sky Compiled Catalogue of 2.5 million stars (Kharchenko 2001) catalogue with known spectral types and which lie within 100 pc (estimated from their Hipparcos parallax measurements) from us. The mean values of $(J-H)$ and $\left(H-K_{\mathrm{s}}\right)$ colours produced are listed in Cols. 3 and 4 of Table 1, respectively. The standard deviations and the number of stars used to evaluate the colours are given in Cols. 5, 6 and 7. The $M_{K}$ values obtained from $M_{V}$ and $(V-K)$ from Cox (2000) are listed in Col. 2. The intrinsic colors of the main sequence stars and giants in the spectral range K8 to M6 were converted from Bessell \& Brett (1988) to the 2MASS system using the transformation equations given by Carpenter (2001). We interpolated the intrinsic colour indices of normal main sequence stars and giants to generate values with a 0.5 spectral sub-type interval.

The locations of stars with various spectral types are identified and marked in Fig. 1. The dashed line is drawn parallel to the Rieke \& Lebofsky (1985) interstellar reddening vector. An unreddened main sequence normal star would move parallel to the reddening vector because of the extinction due to the interstellar dust along the line of sight. But as stars of spectral range A0 to K7 move parallel to the reddening vector, it would be difficult to differentiate between reddened stars in the spectral range A0 to K7 from unreddened M-type normal stars because the M-type stars lie across these vectors. However, for the stars lying below M-type stars loci, it is certain that they are reddened or unreddened normal stars in the spectral range A0 to K7. The stars that occupy the region enclosed by the main sequence star

Table 1. The estimated $(J-H)_{\text {in }}$ and $\left(H-K_{\mathrm{s}}\right)_{\mathrm{in}}$ intrinsic colours for main sequence stars.

| Sp. | $M_{K}$ | $(J-H)_{\text {in }}$ | $\left(H-K_{\mathrm{s}}\right)_{\text {in }}$ | $\sigma_{(J H)_{\mathrm{o}}}$ | $\sigma_{(H K)_{\mathrm{o}}}$ | $N^{\star}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | ---: |
| A0V | 0.65 | -0.026 | 0.029 | 0.030 | 0.028 | 50 |
| A1V | 1.00 | -0.014 | 0.036 | 0.037 | 0.032 | 45 |
| A2V | 1.17 | -0.004 | 0.043 | 0.044 | 0.035 | 43 |
| A3V | 1.35 | 0.020 | 0.048 | 0.049 | 0.032 | 58 |
| A4V | 1.49 | 0.042 | 0.050 | 0.044 | 0.026 | 23 |
| A5V | 1.60 | 0.050 | 0.052 | 0.040 | 0.029 | 24 |
| A6V | 1.62 | 0.058 | 0.051 | 0.037 | 0.024 | 12 |
| A7V | 1.63 | 0.075 | 0.054 | 0.039 | 0.031 | 23 |
| A8V | 1.64 | 0.084 | 0.055 | 0.047 | 0.032 | 8 |
| A9V | 1.72 | 0.090 | 0.059 | 0.033 | 0.026 | 37 |
| F0V | 1.86 | 0.119 | 0.062 | 0.042 | 0.026 | 105 |
| F1V | 2.00 | 0.142 | 0.062 | 0.040 | 0.026 | 9 |
| F2V | 2.15 | 0.162 | 0.064 | 0.035 | 0.025 | 141 |
| F3V | 2.26 | 0.170 | 0.065 | 0.032 | 0.029 | 269 |
| F4V | 2.38 | 0.175 | 0.064 | 0.037 | 0.028 | 28 |
| F5V | 2.49 | 0.195 | 0.068 | 0.036 | 0.027 | 492 |
| F6V | 2.62 | 0.211 | 0.072 | 0.036 | 0.027 | 390 |
| F7V | 2.75 | 0.228 | 0.069 | 0.034 | 0.032 | 387 |
| F8V | 2.88 | 0.235 | 0.069 | 0.038 | 0.029 | 308 |
| F9V | 3.04 | 0.252 | 0.080 | 0.030 | 0.028 | 30 |
| G0V | 3.19 | 0.257 | 0.076 | 0.033 | 0.028 | 396 |
| G1V | 3.24 | 0.265 | 0.078 | 0.035 | 0.028 | 242 |
| G2V | 3.29 | 0.276 | 0.079 | 0.036 | 0.026 | 320 |
| G3V | 3.34 | 0.285 | 0.083 | 0.036 | 0.029 | 477 |
| G4V | 3.45 | 0.302 | 0.084 | 0.032 | 0.031 | 8 |
| G5V | 3.56 | 0.303 | 0.085 | 0.035 | 0.030 | 499 |
| G6V | 3.68 | 0.325 | 0.090 | 0.040 | 0.028 | 188 |
| G8V | 3.90 | 0.347 | 0.093 | 0.048 | 0.030 | 259 |
| G9V | 4.03 | 0.363 | 0.094 | 0.074 | 0.021 | 9 |
| K0V | 4.15 | 0.388 | 0.098 | 0.048 | 0.031 | 207 |
| K5V | 4.35 | 0.53 | 0.604 | 0.160 | 0.031 | 0.035 | 554

$N^{\star}$ the number of stars used to estimate the colours.
loci (i.e., from A0-K7), reddening vector for an A0 type star and below M-type loci were utilized in this work to determine distances to the dark clouds.

### 2.2. Spectral types and $A_{V}$ from near-IR colours

We used a procedure known as "photometric quantification", which is the procedure of inferring spectral type, and hence absolute magnitude and intrinsic colours of normal main sequence stars and giants from NIR photometry. In this method, stars with their $\left(J-K_{\mathrm{s}}\right) \leq 0.75^{1}$ were chosen from the fields containing the cloud. Then a set of dereddened colours [ $(J-H)_{\mathrm{o}}$ and $\left(H-K_{\mathrm{s}}\right)_{\mathrm{o}}$ ] was produced for each star from their observed colours [ $(J-H)$

[^0]and $\left(H-K_{\mathrm{s}}\right)$ ] by using trial values of $A_{V}$ and the Rieke \& Lebofsky (1985) ${ }^{2}$ extinction law in the equations
$(J-H)_{\mathrm{o}}=(J-H)-0.107 \times A_{V}$,
$\left(H-K_{\mathrm{s}}\right)_{\mathrm{o}}=\left(H-K_{\mathrm{s}}\right)-0.063 \times A_{V}$
the trial values of $A_{V}$ were chosen in the range $0-10 \mathrm{mag}$ with an interval of 0.01 mag , though it is evident from Fig. 1 that the maximum extinction that could be traced by this method is limited to $A_{V} \approx 4$ mag. The computed set of dereddened colour indices for a star were then compared with the intrinsic colour indices $\left[(J-H)_{\text {in }} \&\left(H-K_{\mathrm{s}}\right)_{\text {in }}\right.$ ] of normal main sequence stars and giants, produced using the procedure discussed in Sect. 2.1. The best fit of the dereddened colours to the intrinsic colours giving a minimum value of $\chi^{2}$ then yielded the corresponding spectral type and $A_{V}$ for the star. It is known that for a set of $N$ data points $y_{k}$ with associated uncorrelated uncertainties $\sigma_{k}$, the best fit of predicted values $y_{k, \text { mod }}$ (from a model) to these points can be obtained by $\chi^{2}$ minimization defined as (e.g., Gould 2003)
$\chi^{2}=\sum_{k=1}^{N} \frac{\left(y_{k}-y_{k, \bmod }\right)^{2}}{\sigma_{k}^{2}}$.
But while the uncertainties in the magnitudes in $J, H$ and $K_{\mathrm{s}}$ filters could be uncorrelated, the uncertainties in the two colours $(J-H)$ and $\left(H-K_{\mathrm{s}}\right)$ calculated from them with one filter $(H)$ in common would be correlated. In Fig. 2 we show the correlation between the uncertainties in colours of stars from the parts of the sky (L1517, Chamaeleon I, Lupus 3 and NGC 7023) selected in this work. The Spearman rank-order correlation coefficient, $r_{\mathrm{s}}$, calculated for the sources in respective regions indicates a strong correlation between the uncertainties in the $(J-H)$ and $\left(H-K_{\mathrm{s}}\right)$ colours. For correlated uncertainties, Eq. (4) is not valid and one has to use (e.g., Gould 2003)
$\chi^{2}=\sum_{k=1}^{N} \sum_{l=1}^{N}\left(y_{k}-y_{k, \bmod }\right) \boldsymbol{B}_{k l}\left(y_{l}-y_{l, \bmod }\right)$
where $B \equiv C^{-1}$, the inverse of the colour covariance matrix. In the case of the three bands considered here, the colour covariance matrix is given as (Lombardi \& Alves 2001; Maíz-Apellániz 2004)

$C_{J-H, H-K}=\left(\begin{array}{cc}\sigma_{J}^{2}+\sigma_{H}^{2} & -\sigma_{H}^{2} \\ -\sigma_{H}^{2} & \sigma_{H}^{2}+\sigma_{K}^{2}\end{array}\right)$
2 The Rieke \& Lebofsky (1985) law was derived towards the Galactic center in the Arizona-Johnson, not 2MASS, photometric system. But the extinction law derived towards the Galactic center in 2MASS system by Nishiyama et al. (2009), though consistent with the transformed Rieke \& Lebofsky (1985) values, is different from those estimated by Indebetouw et al. (2005) derived towards massive star forming and "field" regions. It was suggested that the extinction law derived towards the Galactic center might be slightly different from those towards other sight lines. Since the maximum extinction that can be traced by our method is limited to $\sim 4 \mathrm{mag}$, the regions studied by us are in close resemblance to those by Indebetouw et al. (2005). However, the ratios, $A_{J} / A_{K}=2.51$ and $A_{H} / A_{K}=1.56$ from the Rieke \& Lebofsky (1985) values are in good agreement with the average values of $A_{J} / A_{K}=2.5 \pm 0.15 \& A_{H} / A_{K}=1.55 \pm 0.08$ obtained by Indebetouw et al. (2005), suggesting that the untransformed values of the Rieke \& Lebofsky (1985) represent the reddening law in the 2MASS system. A similar suggestion was made by Alves et al. (1998) for their observations in the CIT system. We used the values: $A_{J} / A_{V}=0.282$, $A_{H} / A_{V}=0.175, A_{K_{\mathrm{s}}} / A_{V}=0.112$, following Cambrésy et al. (2002), who found that the slope of the reddening vector measured in the 2MASS colour-colour diagram is in better agreement with the Rieke \& Lebofsky (1985) extinction law.


Fig. 2. The relation between the uncertainties in $(J-H)$ and $\left(H-K_{\mathrm{s}}\right)$ colours of the sources in the four parts of the sky studied in this work. The Spearman rank-order correlation coefficient, $r_{\mathrm{s}}$, calculated for the regions are also given. The $r_{\mathrm{s}}$ values indicate a strong correlation between the uncertainties of the two colours. The sources with the uncertainties $\leq 0.035 \mathrm{mag}$ in $J, H$, and $K_{\mathrm{s}}$ filters were considered.
with the off-diagonal elements having non-zero values unlike in the case of uncorrelated uncertainties. The expected value of $\chi^{2}$ after it has been minimized is given by the difference between the number of colours used and the number of parameters to be estimated (e.g., Gould 2003). Because we used two colours, $(J-H)$ and $\left(H-K_{\mathrm{s}}\right)$, to estimate two parameters, spectral type and extinction, the expected values of $\chi^{2}$ should be zero. However, since we used discrete values of $A_{V}(0.01 \mathrm{mag})$ to deredden stars, the dereddened colours of the stars might not match exactly with the intrinsic colours of main sequence and giants to give $\chi^{2}$ values equal to zero but values $\ll 1$. We considered only those solutions which gave a minimized $\chi^{2} \leq 0.1$ for our analysis.

The whole procedure is illustrated in Fig. 1 where we plot the NIR-CC diagram for stars chosen from an arbitrary location on the sky. The arrows are drawn from the observed data points to the corresponding dereddened colours estimated using our method. The maximum extinction values that can be measured using the method are those for A0V type stars ( $\approx 4 \mathrm{mag}$ ). The extinction traced by stars falls as we move towards more late type stars.

### 2.3. Distances to dark molecular clouds

Once the spectral type and the $A_{V}$ of the stars are known, we calculate their distances using the equation
$d(\mathrm{pc})=10^{\left(K_{\mathrm{s}}-M_{K}+5-A_{K}\right) / 5}$.
We excluded the stars classified as giants because of the large uncertainty in their absolute magnitudes. The distance at which the extinction of the stars showed a sudden increase in an $A_{V}$ vs. $d$ plot was considered as the distance to the cloud.

### 2.4. Limitations and uncertainties

The uncertainty in the determination of distances to the clouds depends on the uncertainty in the estimation of distances to the individual stars which is given by
$\sigma_{d}=\sqrt{\left(\sigma_{K_{\mathrm{s}}}^{2}+\sigma_{M_{K}}^{2}+\sigma_{A_{K}}^{2}\right) \times(d / 2.17)^{2}} \quad(\mathrm{pc})$
where $\sigma_{K_{\mathrm{s}}}$ is the photometric uncertainty in $K_{\mathrm{s}}$ band, $\sigma_{M_{K}}$ is the uncertainty in the estimation of the absolute magnitude or the spectral types and $\sigma_{A_{K}}$ is the uncertainty in the $A_{K}(=0.112 \times$ $A_{V}$ ) estimated by the method. The major sources of uncertainty are, however, the photometric errors and the uncertainty due to a different reddening law, these are respectively statistical and systematic in nature.

We allowed a maximum photometric uncertainty ${ }^{3}$ of 0.035 mag for the sources in the $J, H$, and $K_{\mathrm{s}}$ bands. The uncertainty in the absolute magnitudes or the spectral types determined from our method was evaluated by applying it to stars with known distances from Hipparcos taken from the catalogue by Kharchenko (2001). We used stars located within 100 pc of the Sun so that the effect of the extinction could be neglected. Then, Eq. (8) applied to the Hipparcos stars becomes
$\sigma_{d_{\mathrm{H}}}=\sqrt{\left(\sigma_{K_{\mathrm{s}}}^{2}+\sigma_{M_{K}}^{2}\right) \times(d / 2.17)^{2}}$.
In Fig. 3, we show histograms of the difference ( $d_{\mathrm{p}}-d_{\mathrm{h}} \mathrm{pc}$ ) between the distances obtained using spectral types from our method, $d_{\mathrm{p}}(\mathrm{pc})$, and those from the Hipparcos parallax measurements, $d_{\mathrm{h}}(\mathrm{pc})$. Figure 3a shows the distribution when stars with $\sigma_{\pi_{\mathrm{H}}} / \pi_{\mathrm{H}} \leq 0.2$ (ratio of standard error to Hipparcos parallax) from Hipparcos parallax measurements were selected. A Gaussian fit to the distribution shows a $1 \sigma$ deviation of $\sim 18 \mathrm{pc}$. The scatter in the distribution reduced to $\sim 11 \mathrm{pc}$ (Fig. 3b) when stars with $\sigma_{\pi_{\mathrm{H}}} / \pi_{\mathrm{H}} \leq 0.05$ from Hipparcos parallax measurements were selected. The majority ( $85 \%$ ) of the latter stars were found to be located at distances $\lesssim 60 \mathrm{pc}$. We noticed that the mean values of the distributions shown in both Figs. $3 a$ and $b$ are shifted from the expected value of zero. Statistically the error in the mean, given by the standard deviation $/ \sqrt{N}$ ( $N$ is the number of stars here), in both distributions (a) and (b) should be less than 0.36 and 0.47 respectively. The negative shift in (a) could be due to the unaccounted for binary stars which could result in the underestimation of distances using photometry. Adopting a $1 \sigma$ uncertainty of $18 \%$ of the distance for the Hipparcos stars in Eq. (9), the uncertainty in the $\sigma_{M_{K}}$ is estimated to be $\approx 0.4$ which corresponds to an uncertainty in the spectral type of about four subclasses ${ }^{4}$. We assumed 0.4 as the uncertainty in the $\sigma_{M_{K}}$ while calculating the uncertainty in the distances for all the stars.

The $A_{V}$ values obtained for the stars from our method are those which when used to deredden the observed colours (in Eqs. (2) and (3)) gave minimum values of $\chi^{2}$. Therefore the $A_{V}$ values obtained are equivalent to the average value of $A_{V}$ calculated using the individual colours, i.e., $A_{V}=[E(J-H) \times 9.35+$ $\left.E\left(H-K_{\mathrm{s}}\right) \times 15.87\right] / 2$. Since the colours are found to be correlated, the expression for evaluating the uncertainty in $A_{V}$ is
$\sigma\left(A_{V}\right)=\sqrt{4.7^{2} \cdot \sigma_{J H}^{2}+7.9^{2} \cdot \sigma_{H K_{\mathrm{s}}}^{2}+2 \times 37 \times \operatorname{cov}(J H, H K)}$

[^1]

Fig. 3. a) Histogram of the difference between the distances of the main sequence stars from our method and those from the Hipparcos parallax measurements with $\sigma_{\pi_{\mathrm{H}}} / \pi_{\mathrm{H}} \leq 0.2$ are shown with a Gaussian fit to the data over-plotted. b) Same as in a) but for main sequence stars with $\sigma_{\pi_{\mathrm{H}}} / \pi_{\mathrm{H}} \leq 0.05$.
where $\sigma_{J H}^{2}=\sigma_{J}^{2}+\sigma_{\mathrm{H}}^{2}, \sigma_{H K_{\mathrm{s}}}^{2}=\sigma_{H}^{2}+\sigma_{K_{\mathrm{s}}}^{2}$ and $\operatorname{cov}(J H, H K)=$ $r_{\mathrm{s}} \times \sigma_{J H} \sigma_{H K_{\mathrm{s}}}$. Using the maximum photometric uncertainty of 0.035 mag allowed in the individual bands and the highest value of 0.93 for $r_{\mathrm{s}}$ (obtained towards L1517), the maximum uncertainty in the $A_{V}$ was found to be $\sim 0.6 \mathrm{mag}$. The uncertainty in the distances to the individual stars, Eq. (8), is therefore dominated by the uncertainty in the determination of the spectral types or $M_{K}$ values.

The true spectral type and the $A_{V}$ of the stars evaluated from our method directly depend on the reddening or extinction law that we adopt. Unlike in visual and UV wavelengths where the extinction law is known to vary along different sight lines in the Galaxy (Mathis 1990), the similarity in the wavelength dependence of extinction for very different line of sights passing through different environments like the giant H II region (RCW 49) and the "field" region (Indebetouw et al. 2005) suggests that the extinction law is almost universal in the near-IR regime considered here. Consequently, the observed colours can be corrected for the reddening to determine the spectral type of the stars using a standard extinction law. But any deviation of the reddening law from the standard extinction law towards the


Fig. 4. Procedure used to determine distances to the clouds from an $A_{V}$ vs. $d$ as described in Sect. 2.5. The extinction of stars in situations with and without the presence of a cloud in the field-of-view of a region is represented by open and closed circles. The mean value of extinction in each distance bin in case of with and without cloud is shown by red and blue solid lines, respectively.
region containing the clouds could introduce a systematic error in the distances estimated using our method.

### 2.5. Determination of cloud distance from $A_{V}$-distance plot

The presence of an interstellar cloud should reveal itself by a sudden onset of extinction at a certain distance. If the cloud is isolated and is the first interstellar feature encountered, the excess should almost equal to zero in front of the cloud (e.g., Straižys et al. 1992; Whittet et al. 1997; Knude \& Høg 1998; Alves \& Franco 2006; Lombardi et al. 2008). In this technique, the distance to the cloud is typically estimated from the first star that shows a significant reddening. The estimated distance is thus regarded as an upper limit due to the lack of information on the closeness of this background star to the cloud. In our method, because we used a relatively large number of stars projected towards the cloud, we often found an almost continuous distribution of stars with high extinction beyond a certain distance.

In Fig. 4 we show a schematic figure of an $A_{V}$ vs. $d$ plot for an ideal situation along a line of sight through the Galactic plane. The extinction suffered by the stars increases as a function of their distance (typically, $\sim 1.8 \mathrm{mag} \mathrm{kpc}^{-1}$ is the value considered for the standard mean rate of diffuse-ISM extinction in the Galactic plane near the Sun, Whittet 1992) is represented by filled circles. If we group the stars in distance bins and calculate the mean $A_{V}$ of the stars for each bin, they would follow the blue curve as shown in Fig. 4. The open circles represent the situation when a cloud of constant extinction is introduced at a certain distance in the field-of-view. Now, the mean $A_{V}$ of the stars in the bins behind the cloud would show a jump equal to the extinction introduced by the cloud (red curve). The distance at which the first sudden increase in the mean $A_{V}$ occurs could be considered as the distance to the cloud (marked by the vertical dashed line in Fig. 4).

In real cases, we grouped the stars into distance bins of bin width $=0.09 \times$ distance. The centers of each bin were separated by the half of the bin width. Because there exist very few stars at smaller distances, the mean of the distances and the $A_{V}$
of the stars in each bin were calculated by taking 1000 pc as the initial point and proceeding towards smaller distances (see Fig. 17). The mean distance of the stars in the bin at which a significant drop in the mean of the extinction occurred was taken as the distance to the cloud and the average of the uncertainty in the distances of the stars in that bin was taken as the final uncertainty in our distance determination for the clouds. For all the clouds studied in this work, the vertical dashed lines in $A_{V}$ vs. $d$ plots, used to mark the cloud distances, were drawn at distances deduced from the above procedure (see Fig. 17). The error in the mean values of $A_{V}$ were calculated using the expression standard deviation $/ \sqrt{N}$ where $N$ is the number of stars in each bin.

## 3. The data

We extracted $J, H$ and $K_{\mathrm{s}}$ magnitudes of stars from the 2MASS All-Sky Catalog of Point Sources (Cutri et al. 2003) that satisfied the following criteria,

1. photometric uncertainty $\sigma \leq 0.035$ in all the three filters;
2. photometric quality flag of "AAA" in all the three filters, i.e., signal-to-noise ratio $(S N R)>10$.

The maximum allowed uncertainty of 0.035 magnitude in the photometry was chosen primarily as a trade-off between the minimum scatter in the data points and the availability of enough stars in the background and the foreground of the clouds to bracket the cloud distances. One added advantage of choosing the data with high quality was that the limiting magnitude becomes brighter and therefore, the selected stars would sample the lower density regions of the clouds where a normal reddening law is believed to be the most valid (Straižys et al. 1982; Kandori et al. 2003).

## 4. Distances to the molecular clouds L1517, Chamaeleon I, Lupus 3 and NGC 7023

We applied our method to four clouds, L1517, Chamaeleon I (Cha I), Lupus 3 (Lup 3) and NGC 7023, to determine distances to them. We preferred these four regions as they are structurally less complex and have had their distances estimated previously. Also, these clouds harbour stars (inferred from the associated nebulosity) for which distances from the Hipparcos parallax measurements are available (van den Ancker et al. 1998; Bertout et al. 1999), which would help constrain the parent cloud distances better. In Table 2, we list the distances to them compiled from the literature. The central coordinates, total number of stars selected and the number of stars classified as dwarfs and that were used to determine distances to the individual clouds are listed in Cols. 2-5 of Table 3 respectively.

### 4.1. L 1517

In Fig. 5 we present the $2^{\circ} \times 2^{\circ}$ high resolution extinction map ${ }^{5}$ produced by Dobashi et al. (2005) of the region containing L1517. The contours are drawn at $0.5,1.0,1.5$ and 2 mag levels. The Herbig Ae star, AB Aur (Herbig 1960; The et al. 1994), associated with the reflection nebulosity, GN 04.52.5.02 (Magakian 2003), is identified in Fig. 5.

[^2]Table 2. Distances to the clouds L1517, Cha I, Lup 3 and NGC 7023 compiled from the literature.

| Cloud <br> Identification | Distances quoted <br> in the literature (pc) | Associated <br> star | Ref. |
| :--- | :--- | :--- | :--- |
| L 1517 | $144_{-17}^{+23+}$ | AB Aur | 1 |
| Chamaeleon 1 | $115-215,158$ | HD 97048 | 2,3 |
|  | $150 \pm 15,175_{-20}^{+27}$ |  | 4,1 |
| Lupus 3 | $130-170,140$ | HR 5999 | $5,6,7$ |
|  | $100,208_{-32}^{+46,140 \pm 2}$ |  | $8,9,10$ |
| NGC 7023 | $440,429_{-90}^{+199}, 288 \pm 25$ | HD 200775 | 11,1 |
|  | $288 \pm 25$ |  | 12 |

1-Bertout et al. (1999); 2-Schwartz (1991); 3-Franco (1991); 4-Whittet et al. (1997); 5-Murphy et al. (1986); 6-Krautter (1991); 7-Hughes et al. (1993); 8-Knude \& Høg (1998); 9-van den Ancker et al. (1998); 10-de Zeeuw et al. (1999); 11-Viotti (1969); 12-Straižys et al. (1992).
${ }^{\dagger}$ The Hipparcos distance to the AB Aur.

Table 3. The central coordinates, the number of stars selected from each field and the stars classified as dwarfs.

| Field <br> Id. | $\alpha(J 2000)$ <br> $\left({ }^{\circ}\right)$ | $\delta(J 2000)$ <br> $\left({ }^{\circ}\right)$ | Total <br> stars | Stars classified <br> as dwarfs |
| :--- | :---: | :---: | :---: | ---: |
| L1517 |  |  |  |  |
| F1 | 74.166973 | +30.557486 | 395 | 292 |
| F2 | 73.670855 | +31.053637 | 339 | 260 |
| F3 | 73.555237 | +30.138617 | 508 | 402 |
| F4 | 73.080485 | +30.648516 | 299 | 230 |
| Total |  |  | 1541 | 1184 |
| Chamaeleon I |  |  |  |  |
| F1 | 169.399281 | -76.719727 | 521 | 360 |
| F2 | 165.432583 | -76.312943 | 494 | 352 |
| F3 | 167.607235 | -77.711487 | 235 | 161 |
| F4 | 163.267247 | -77.305313 | 360 | 260 |
| F5 | 165.209741 | -78.675217 | 502 | 351 |
| F6 | 160.921746 | -78.242210 | 550 | 389 |
| Total |  |  | 2662 | 1873 |
| Lupus 3 |  |  |  |  |
| F1 | 242.405769 | -39.179062 | 490 | 335 |
| F2 | 241.008065 | -39.471130 | 814 | 561 |
| F3 | 243.129521 | -38.354233 | 487 | 308 |
| Total |  |  | 1791 | 1204 |
| NGC 7023 |  |  |  |  |
| F1 | 315.431610 | +68.160049 | 310 | 230 |
| F2 | 314.744713 | +67.160805 | 394 | 281 |

Some of the stars classified as main sequence by our method could actually be giants. Such erroneous classifications could lead to an underestimation of their distances which could create confusion with the real increase in the extinction due to the presence of a cloud. This problem could be circumvented if we divide a large field, containing the cloud, into small sub-fields. While the rise in the extinction due to the presence of a cloud should occur almost at the same distance in all the fields, if the whole cloud is located at the same distance, the wrongly classified stars in the sub-fields would show high extinction not at the same but at random distances. We divided the region containing L1517 into four fields as shown in Fig. 5. Each field is of $40^{\prime} \times 40^{\prime}$ in area. The central coordinates, the number of stars selected after applying all the selection criteria ( $\sigma \leq 0.035$, $S N R>10$, and $\left.\left(J-K_{\mathrm{s}}\right) \leq 0.75\right)$ and the number of stars


Fig. 5. The $2^{\circ} \times 2^{\circ}$ extinction map produced by Dobashi et al. (2005) containing L1517 is shown with the fields F1-F4, each covering $40^{\prime} \times$ $40^{\prime}$ area, marked and labeled. The contours are drawn at 0.5, 1.0, 1.5 and 2.0 mag levels. The stars used for estimating distance to L1517 are represented by filled circles. The position of the Herbig Ae star, AB Aur, is also identified and labeled. The star with high extinction ( $A_{V}>1 \mathrm{mag}$ ) at a distance closer than the estimated distance of the cloud is identified using the star symbols.
classified as dwarfs by our method from each field are listed in Table 3. In Fig. 6 we show the NIR-CC diagrams for the stars selected from F1-F4. The stars with their photometric errors $\leq 0.035 \mathrm{mag}$ in $J, H$, and $K_{\mathrm{s}}$ bands are shown as dots and among them the stars with $\left(J-K_{\mathrm{s}}\right) \leq 0.75$ are identified by open circles. The stars classified as dwarfs using our method are identified by filled circles (in white) in Fig. 5. It can be noticed that very few sources are selected towards the high extinction regions due to our data selection criteria discussed in Sect. 3.

In Fig. 7 we show the $A_{V}$ vs. $d$ plot for the stars from all the fields combined. The solid curve shows the increase in the extinction towards the Galactic latitude $b=-7.98^{\circ}$ as a function of distance, produced using the expressions given by Bahcall \& Soneira (1980, BS80, hereafter). Clearly, the stars with high extinction located beyond 167 pc , marked by the vertical dashed line (drawn using the procedure described in the Sect. 2.5 , see Fig. 17), are distributed almost continuously. There exists one star with $A_{V}>1$ mag at a distance closer than 167 pc . The $A_{V}$ vs. $d$ plots for the stars from the individual fields F1-F4 are shown in Fig. 8. The stars showing high extinction at 167 pc are located in the fields F1 and F4, which contain the major parts of the cloud. The high extinction star closer than 167 pc in the field F4 is identified using filled star symbol in Fig. 5. Because the star is projected within the cloud boundary, it is likely to be a misclassified source located behind the cloud.

The distances to the molecular clouds in the Taurus-Auriga complex is considered as 140 pc based on the star count, photometric distances of reflecting nebulae and reddening versus distance diagrams of the field stars towards the entire complex. But the most reliable estimation of the distance to L1517 is based on the Hipparcos parallax measurements of the two stars, AB Aur and SU Aur, found associated with the cloud. The Hipparcos


Fig. 6. The NIR-CC diagrams for the stars selected from the four fields containing L1517 are shown. The dots represent all the stars in a given field with photometric errors $\leq 0.035$ mag in $J, H$, and $K_{\mathrm{s}}$ bands and open circles represent stars with their $\left(J-K_{\mathrm{s}}\right) \leq 0.75$.


Fig. 7. The $A_{V}$ vs. $d$ plot for all the stars obtained from the fields F1-F4 combined towards L1517. The dashed vertical line is drawn at 167 pc inferred from the procedure described in the Sect. 2.5. The solid curve represents the increase in the extinction towards the Galactic latitude of $b=-7.98^{\circ}$ as a function of distance produced from the expressions given by Bahcall \& Soneira (1980). The error bars are not shown on all the stars for better clarity.
distances to these stars are $144_{-17}^{+23}$ and $152_{-34}^{+63}$ respectively. The distance of $167 \pm 30 \mathrm{pc}$ to L 1517 determined from our method are found to be in good agreement with the distances of AB Aur and SU Aur within the uncertainties.

### 4.2. Chamaeleon 1

The region containing the cloud Cha I is shown in the $5^{\circ} \times 5^{\circ}$ extinction map produced by Dobashi et al. (2005) in Fig. 9. The


Fig. 8. The $A_{V}$ vs. $d$ plots for the stars from the fields F1-F4 towards L1517. The dashed vertical line is drawn at 167 pc inferred from the procedure described in Sect. 2.5. The solid curve represents the increase in the extinction towards the Galactic latitude of $b=-7.98^{\circ}$ as a function of distance produced from the expressions given by Bahcall \& Soneira (1980).
contours are drawn at $0.5,1.0,1.5$ and 2.0 mag levels. Since the cloud is extended over a region of about $3^{\circ} \times 2^{\circ}$ area, we divided the whole region into six sub-fields of $1^{\circ} \times 1^{\circ}$ each as identified in Fig. 9. The central coordinates, total number of stars selected after applying all the selection criteria ( $\sigma \leq 0.035, S N R>10$, and $\left.\left(J-K_{\mathrm{s}}\right) \leq 0.75\right)$ and the number of stars classified as dwarfs using our method in each field are listed in Table 3. Stars classified as dwarfs are identified in Fig. 9 using filled circles. Here too we find very few sources towards the high extinction regions.

In Fig. 10, we show the $A_{V}$ vs. $d$ plot for stars from all the fields combined. The solid curve represents the increase in the extinction towards the Galactic latitude of $b=-15.9^{\circ}$ as a function of distance produced using the expressions given by BS80. Using the procedure described in Sect. 2.5, we found a drop in the extinction at 151 pc (marked by the vertical dashed line in Fig. 10) as shown in Fig. 17. But the drop in the extinction at 151 pc was not as conspicuous as was observed towards L1517. Three stars with $A_{V} \geq 1$ were found to be located closer than 151 pc (identified in Fig. 9 using filled star symbols). We noticed another drop in the extinction at 128 pc (see Fig. 17). In


Fig. 9. The $5^{\circ} \times 5^{\circ}$ extinction map produced by Dobashi et al. (2005) containing Cha I is shown with the contours drawn at $0.5,1.0,1.5$ and 2.0 mag levels. The fields F1-F6 selected towards Cha I are marked and labeled in the figure. The stars classified as dwarfs using our method are identified with filled circles. Each field is of $1^{\circ} \times 1^{\circ}$ in area. The locations of the stars HD 97300 and HD 97048 are identified and labeled. The stars with high extinction $\left(A_{V}>1 \mathrm{mag}\right)$ at distances closer than the estimated distance of the cloud are identified using the star symbols.


Fig. 10. The $A_{V}$ vs. $d$ plot for the stars from all the fields (F1-F6) of Cha I combined. The vertical dashed line is drawn at 151 pc inferred from the procedure described in Sect. 2.5. The vertical dotted line indicates the distance of a possible foreground dust layer at $\sim 128 \mathrm{pc}$. The solid curve represents the increase in the extinction towards the Galactic latitude of $b=-15.9^{\circ}$ as a function of distance produced from the expressions given by BS80. For clarity, the error bars are not shown on all the stars.

Fig. 11, we show the $A_{V}$ vs. $d$ plot for stars in the individual fields. The dotted line is drawn at 128 pc . The stars with extinction significantly above the values expected from the BS80 model seem to be present in all the fields at or beyond 151 pc , unlike the distribution of the stars closer than 151 pc , implying


Fig. 11. The $A_{V}$ vs. $d$ plot for stars in the fields F1-F6 selected towards Cha I. The dashed vertical line is at 151 pc inferred from the procedure described in Sect. 2.5. The dotted vertical line is drawn at 128 pc . The solid curve represents the increase in the extinction towards the Galactic latitude of $b=-15.9^{\circ}$ as a function of distance produced from the expressions given by Bahcall \& Soneira (1980).
that the foreground dust layer at 128 pc is very patchy in distribution.

Earlier attempts to determine distances to Cha I were made by Whittet et al. (1987), Franco (1991), Schwartz (1991) and Whittet et al. (1997). Schwartz (1991) reported a value of


Fig. 12. The $4^{\circ} \times 4^{\circ}$ extinction map produced by Dobashi et al. (2005) containing Lup 3 with the contours drawn at $0.5,1.0,1.5,2.0,2.5$ and 3.0 mag. The three fields F1-F3 each of $1^{\circ} \times 1^{\circ}$ area containing the cloud and the Herbig Ae/Be star, HR 5999, are identified and labeled in the figure. The stars with high extinction $\left(A_{V}>1 \mathrm{mag}\right)$ at distances closer than the estimated distance of the cloud are identified using star symbols.

115-215 pc, Franco (1991) assigned 158 pc while Whittet et al. (1997) deduced a most probable distance of $150 \pm 15 \mathrm{pc}$ based on the reddening vs. distance plot for stars towards Cha I. Independent distance estimates based on parallax measurements made by the Hipparcos satellite are available for two stars, HD 97048 and HD 97300 (van den Ancker et al. 1998; Bertout et al. 1999), which illuminate prominent reflection nebulae, Ced 111 and Ced 112 respectively. The parallax distances of HD 97048 and HD 97300 were estimated to be $175_{-20}^{+27}$ and $188_{-30}^{+44}$ respectively (Bertout et al. 1999). But the spectrophotometric distances to these stars were estimated to be $180 \pm 20$ and $152 \pm 18$ (Whittet et al. 1997). These stars are located in F3 and F1 as identified in Fig. 9. Based on $u v b y \beta$ photometry of 1017 stars, Corradi et al. (1997) showed the presence of a dust layer which is a part of a large scale distribution extended over a large area covering Chamaeleon, Musca and Coalsack at a distance of $\sim 150 \mathrm{pc}$. In a subsequent work (Corradi et al. 2004) they showed the presence of gas as close as 60 pc from the Sun using interstellar Na I D absorption lines towards 63 B-type stars. The dust layer found in Figs. 10 and 11 at 128 pc could be a part of the forground dust layer shown by Corradi et al. (1997, 2004). The onset of extinction in $A_{V}$ vs. $d$ plots for both F1 and F3 are found to occur consistently at $\sim 151 \mathrm{pc}$ (Fig. 11). Based on our method, the present determination of $151 \pm 28 \mathrm{pc}$ for Cha 1 is found to be in good agreement with the most probable distance estimated for the cloud by Whittet et al. (1997) and Corradi et al. (1997, 2004).

### 4.3. Lupus 3

The region containing the Lup 3 cloud is shown in the $4^{\circ} \times 4^{\circ}$ extinction map produced by Dobashi et al. (2005) in Fig. 12. The contours are drawn at 0.5 to 3.0 with an interval of 0.5 mag . The three fields containing the cloud are identified with square boxes.


Fig. 13. The $A_{V}$ vs. $d$ plot for the stars from all the fields F1-F3 of Lup 3 combined. The dashed at 157 pc inferred from the procedure described in Sect. 2.5. The error bars are not shown on all the stars for clarity.

The field, F1, contains the Herbig Ae star, HR 5999 (The et al. 1994). The central coordinates, total number of stars selected after applying all the selection criteria ( $\sigma \leq 0.035, S N R>10$, and $\left.\left(J-K_{\mathrm{s}}\right) \leq 0.75\right)$ and the number of stars classified as dwarfs (identified in Fig. 12 with filled circles) are listed in Table 3. As seen in the cases of L1517 and Cha I, towards Lupus 3 too, we found very few stars towards the dense parts of the cloud.

In Fig. 13, we show the $A_{V}$ vs. $d$ plot for the stars from all three fields combined. The solid curve represents the increase in the extinction towards the Galactic latitude of $b=+9.4^{\circ}$ as a function of distance produced using the expressions given by BS80. We found a sharp drop in the extinction at 157 pc , as shown in Fig. 17, estimated using the procedure described in Sect. 2.5 (marked with the vertical dashed line in Fig. 13). There are three stars with $A_{V} \geq 1$ at distances closer than 157 pc in Fig. 13. In Fig. 14, we show the $A_{V}$ vs. $d$ plot for the individual fields towards Lup 3. The source located at 49 pc and $A_{V}=0.5$ in F2 is identified as HD 143 261, a giant, classified as K1/K2III (Simbad database). Of the two stars with $A_{V} \geq 1$ in F3, the one located at 110 pc and $A_{V}=1.9$ is identified as HD 145355 classified as A1III (Simbad database).

Most of the earlier studies have determined or assumed distances in the range $130-170 \mathrm{pc}$ either by considering Lupus to be associated with the Scorpius-Centaurus association (Murphy et al. 1986; Krautter 1991) or from spectroscopic parallaxes of early-type stars thought to be members of Scorpius-Centaurus association (Hughes et al. 1993). But Knude \& Høg (1998) based on extinction of stars located in the foreground and the background of Lupus complex with their distances estimated from Hipparcos parallaxes, determined a closer distance of 100 pc to Lup 3. However, the Hipparcos parallax distance of HR 5999, illuminating the nebula GN 16.05.2 (Magakian 2003), is estimated to be $208_{-32}^{+46} \mathrm{pc}$ (van den Ancker et al. 1998) farther than most of the distance estimates. The best distance estimation currently available for the Centaurus-Lupus subgroup of the ScorpiusCentaurus association is $140 \pm 2 \mathrm{pc}$ derived by de Zeeuw et al. (1999) based on Hipparcos parallaxes. Using our method, we determined a distance of $157 \pm 28 \mathrm{pc}$ to the Lup 3 cloud.

### 4.4. NGC 7023

The reflection nebula NGC 7023 is illuminated by the Herbig Be star HD 200775 (The et al. 1994). The $3^{\circ} \times 3^{\circ}$ extinction map


Fig. 14. The $A_{V}$ vs. $d$ plots for stars from the individual fields F1-F3 of Lup 3. The dashed line is at 157 pc inferred from the procedure described in the Sect. 2.5. The dot-dashed line shows the increase in extinction towards the Galactic latitude of $b=+9.4^{\circ}$ as a function of distance estimated using expressions given by BS80. The source marked with an open triangle is identified as HD 143 261, a giant (K1/K2III).
of the region containing NGC 7023 is shown in Fig. 15. The contours are drawn at $0.6,1.0,1.5,2.0,2.5$ and 3.0 mag. Two fields of $1^{\circ} \times 1^{\circ}$ were chosen towards the region. The central coordinates, total number of stars selected after applying all the selection criteria ( $\sigma \leq 0.035, S N R>10$, and $\left(J-K_{\mathrm{s}}\right) \leq 0.75$ ) and the number of stars classified as dwarfs are listed in Table 3. These stars are identified with filled circles in Fig. 15.

In Fig. 16, we present the $A_{V}$ vs. $d$ plot for stars from the fields F1 and F2. The solid curve shows the increase in the extinction towards the Galactic latitude of $b=+14.1^{\circ}$ as a function of distance produced using the expressions given by BS80. Clearly, a drop in the extinction is noticeable at two distances for F 1 , at $\approx 310 \mathrm{pc}$ and at $\approx 410 \mathrm{pc}$. In F2, though there exists a steplike appearance again at distances $\approx 310$ and $\approx 410 \mathrm{pc}$, a number of stars show relatively high extinction at a much closer distance of $\approx 200 \mathrm{pc}$, a distance where an increase in the $A_{V}$ of $\geq 1.5 \mathrm{mag}$ also seen in F1 but for only two stars. From the $A_{V}$ vs. $d$ plot for stars in F1 and F2, we infer the presence of at least three layers of dust grains along the line of sight towards NGC 7023. While the dust components at $\approx 310 \mathrm{pc}$ and $\approx 410 \mathrm{pc}$ are dominant towards F1, the components at $\approx 200 \mathrm{pc}$ and $\approx 310 \mathrm{pc}$ are dominant towards F2. Using the procedure discussed in Sect. 2.5, we found a significant drop in the mean value of $A_{V}$ at 408 pc (the distance
G. Maheswar et al.: NIR-photometric distances to molecular clouds


Fig. 15. The $3^{\circ} \times 3^{\circ}$ extinction map of the region containing NGC 7023 is shown with the fields F1 and F2 ( of $1^{\circ} \times 1^{\circ}$ each) identified. The contours are drawn at $0.6,1.0,1.5,2.0,2.5$ and 3.0 mag. The arrow shows the location of the Herbig Ae/Be star HD 200775. The stars (only from F1) with high extinction ( $A_{V}>1 \mathrm{mag}$ ) at distances closer than the estimated distance of the cloud is identified using star symbols.


Fig. 16. The $A_{V}$ vs. $d$ plots for the stars in F1 and F2 towards NGC 7023. The dashed lines is drawn at 408 pc inferred from the the procedure discussed in Sect. 2.5. The dotted lines indicate the locations of two additional dust layers possibly present at $\approx 200 \mathrm{pc}$ and $\approx 310 \mathrm{pc}$. The solid curve represents the increase in extinction towards the Galactic latitude of $b=+14.1^{\circ}$ as a function of distance estimated using expressions given by BS80. The error bars are not shown on all the stars for clarity.


Fig. 17. The mean values of $A_{V}$ vs. the mean values of distance plot for L1517, Cha I, Lupus 3 and NGC 7023 produced using the the procedure discussed in the Sect. 2.5 to determine distances to the clouds. The error bars on the mean $A_{V}$ values were calculated using the expression, standard deviation $/ \sqrt{N}$, where $N$ is the number of stars in each bin.
marked by the vertical dashed line). We used the stars from F1 alone as the situation is much more complex towards F2.

The most widely quoted distances to LDN 1167/1174 or the NGC 7023 and LDN 1147/1158 (clouds located $\sim 2^{\circ}$ west of NGC 7023) groups are $288 \pm 25$ and $325 \pm 13$ pc respectively (Straižys et al. 1992). They used the Vilnius photometric system to classify the stars and also to get the interstellar reddening. They obtained a distance of $325 \pm 13 \mathrm{pc}$ to LDN $1147 / 1158$ by taking an average of 10 stars showing $A_{V} \geq 0.45$ which were distributed in the range 240-380 pc. For LDN 1167/1174 or NGC 7023 group, they assigned a distance of $288 \pm 25 \mathrm{pc}$ again by taking the average distance of 4 considerably reddened stars. They preferred a distance of 275 pc to HD 200775 by assuming it to be a B3Ve star.

The Hipparcos parallax distance to HD 200775 is estimated to be $429_{-90}^{+156} \mathrm{pc}$ (van den Ancker at al. 1998; Bertout et al. 1999). The sharp rise in the extinction in Fig. 16 at 408 pc could be due to the association of the cloud with HD 200775 and the sharp
rise in $A_{V}$ at 305 pc could be due to a foreground dust component. Kun (1998) showed the presence of dust components at three characteristic distances: 200, 300, and 450 pc , based on a cumulative distribution of field star distance moduli in Wolf diagrams. We found an additional component of dust at 170 pc towards the western parts of NGC 7023, especially towards the LDN 1171 and LDN 1147/1158 groups. The results will be presented in a forthcoming paper. Using a total of 230 stars classified as dwarfs, we determined a distance of $408 \pm 76 \mathrm{pc}$ to the NGC 7023 cloud.

## 5. Discussion

The current determination of distances to L1517, Cha I, Lupus 3 and NGC 7023 is based on the NIR photometry of a relatively large number of stars, 1184, 1873, 1204 and 230, respectively (Table 3), projected onto them. Distances to these clouds estimated using our method are found to be in good agreement with most of the previous estimations within the error (see Table 2).

Based on our method and the NIR data from 2MASS, one could determine distances to any cloud located within $\sim 500 \mathrm{pc}$ from the Sun (depending upon the Galactic latitude of the clouds). However, the complex nature of the interstellar medium is clearly evident in the $A_{V}$ vs. $d$ plots of Cha 1 and NGC 7023, as shown in Figs. 10 and 16, respectively. There exist, clearly, at least two layers of dust components in the direction towards these two regions. If the first layer of the dust component happens to be denser but uniform in distribution, it would reduce the probability of finding the second dust layer. However, if the first layer is denser but patchy in distribution, we would find sources with high extinction at shorter distances before we would notice a wall of stars with high extinction due to the second layer of dust component. In that case, differentiating stars with spurious spectral classification from the actual rise in extinction due to the presence of dust components would become difficult. Dividing a larger field into sub-fields could help us decipher the $A_{V}$ vs. $d$ plots of a cloud. It is indeed noted that the presence of dominant dust components is revealed in $A_{V}$ vs. $d$ plots for different fields towards Cha 1 (Fig. 10) and NGC 7023 (Fig. 16) where the effects are most conspicuous.

The size of the sub-fields required towards a cloud should be decided on the basis of its galactic location. We noticed that in a given field, our method consistently assigns $\approx 70 \%$ of the total selected stars as dwarfs, after applying all the selection criteria ( $\sigma \leq 0.035, S N R>10$, and $\left(J-K_{\mathrm{s}}\right) \leq 0.75$ ). Our experience from the application of the method to the four clouds in the previous sections showed that we require at least 100200 dwarfs to discern the presence of a sudden rise (or drop) in the extinction in the $A_{V}$ vs. $d$ plot of a field. The number of stars decreases with Galactic latitude, thus requiring a larger field as we go towards relatively high galactic latitudes. At higher galactic latitudes, however, the confusion due to multiple dust components in the line of sight would be reduced as the majority of the clouds tend to occupy regions closer to the Galactic plane.

For small and isolated clouds, it would be difficult to divide the field containing the cloud into sub-fields with a sufficient number of stars projected on to them and therefore the determination of distance is made correspondingly hard. In such cases, clouds that are spatially closer and show similar radial velocities should be chosen. Clouds located in the close proximity and with similar radial velocities are believed to be at similar distances.

## 6. Conclusions

We present a method to determine distances to molecular clouds using the prodigious amount of NIR data provided by the 2MASS. The method involves the following steps:

1. Extract $J, H, K_{\mathrm{s}}$ magnitudes with photometric uncertainty $\leq 0.035$ and $S N R>10$ of stars projected onto the fields containing the cloud from the 2 MASS database.
2. Select stars with their $\left(J-K_{\mathrm{s}}\right) \leq 0.75$ to eliminate M-type stars from the analysis, as unreddened M-type stars located across the reddening vectors of A0-K7 dwarfs make it difficult to differentiate the reddened A0-K7 dwarfs from the unreddened M-type stars.
3. A set of dereddened colours for each star is produced from their observed colours by using a range of trial values of $A_{V}$ ( $0-10 \mathrm{mag}$ ) and the Rieke \& Lebofsky (1985) extinction law in Eqs. (2) and (3).
4. Computed sets of dereddened colours of a star are then compared with the intrinsic colours of the normal main sequence stars. The intrinsic colours of the main sequence stars in the spectral range A0-K7 are produced from the 2MASS data of stars with known spectral types and those located within 100 pc , calculated from the Hipparcos parallaxes. The best match giving a minimum value of $\chi^{2}$ as defined in Eq. (5) then yields the spectral type and $A_{V}$ corresponding to that intrinsic colour.
5. Once the spectral types and $A_{V}$ values of the stars are known, their distances are estimated using the distance Eq. (1). Only those stars that are classified as dwarfs are considered for the determination of distances, as the absolute magnitudes of giants are highly uncertain.
We applied the method first to main sequence stars located at distances $\leq 100 \mathrm{pc}$ (from Hipparcos parallaxes) and have their spectral types known. We showed that our method can determine the spectral types of the stars with a $1 \sigma$ uncertainty of about four subclasses which propagates into the uncertainty in the distances of the stars by $\leq 20 \%$. By applying the method to four clouds, L1517, Chamaeleon I, Lupus 3 and NGC 7023, which harbour nebulous stars with Hipparcos parallax distances known, we determined a distance of $167 \pm 30,151 \pm 28,157 \pm 29$ and $408 \pm$ 76 pc , respectively. These are found to be in good agreement with the most accurate distances available for them in the literature. We found additional components of dust layers at $\approx 100 \mathrm{pc}$ towards Chamaeleon I and at $\approx 200$ and $\approx 310$ pc towards NGC 7023. In the future, we will extend the application of this method to many more molecular clouds to determine their distances using the 2MASS NIR photometry.
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[^0]:    ${ }^{1}$ This criterion allows us to avoid M-type stars from the analysis. Also, the classical T Tauri stars (CTTSs), often found to be associated the molecular clouds, occupy a well defined locus in the NIR-CC diagram (Meyer et al. 1997) as shown in Fig. 1 which intercepts with the main sequence loci at $\left(J-K_{\mathrm{s}}\right) \approx 0.75$. Thus the criterion of $\left(J-K_{\mathrm{s}}\right) \leq 0.75$ would enable us to reject most of the CTTSs as well.

[^1]:    ${ }^{3}$ The photometric errors considered in this work in the $J, H$, and $K_{\mathrm{s}}$ magnitudes from the 2MASS database include the corrected band photometric uncertainty, nightly photometric zero point uncertainty, and flat-fielding residual errors (Cutri et al. 2003).
    ${ }^{4}$ The change in $M_{K}$ with respect to the spectral types in the range A0V-K7V could be fitted with a function, $M_{K} \approx 0.1 \times \mathrm{Sp}$. type +0.98 .

[^2]:    5 The extinction map, covering the entire region in the galactic latitude range $|b| \leq 40^{\circ}$ derived using the optical database "Digitized Sky Survey I" and the traditional star-count technique, was produced in two angular resolutions of $6^{\prime}$ and $18^{\prime}$ (Dobashi et al. 2005).

